



19 September 2023

Kilonovae: the cosmic foundries of heavy elements (Part I)

Elena Pian
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Gravitational waves are "ripples" of Spacetime, produced by accelerated masses and predicted by General Relativity theory (Einstein 1916, 1918)

Über Gravitationswellen.

VON A. EINSTEIN.

(Vorgelegt am 31. Januar 1918 [s. oben S. 79].)

Die wichtige Frage, wie die Ausbreitung der Gravitationsfelder erfolgt, ist schon vor anderthalb Jahren in einer Akademienarbeit von mir behandelt worden¹. Da aber meine damalige Darstellung des Gegenstandes nicht genügend durchsichtig und außerdem durch einen bedauerlichen Rechenfehler verunstaltet ist, muß ich hier nochmals auf die Angelegenheit zurückkommen.

Wie damals beschränke ich mich auch hier auf den Fall, daß das betrachtete zeiträumliche Kontinuum sich von einem galileischen nur sehr wenig unterscheidet. Um für alle Indizes

$$g_{\alpha\beta} = -\delta_{\alpha\beta} + \gamma_{\alpha\beta} \quad (1)$$

setzen zu können, wählen wir, wie es in der speziellen Relativitätstheorie üblich ist, die Zeitvariable x_4 rein imaginär, indem wir

$$x_4 = it$$

setzen, wobei t die »Lichtzeit« bedeutet. In (1) ist $\delta_{\alpha\beta} = t$ bzw. $\delta_{\alpha\beta} = 0$, je nachdem $\alpha = \beta$ oder $\alpha \neq \beta$ ist. Die $\gamma_{\alpha\beta}$ sind gegen t kleine Größen, welche die Abweichung des Kontinuums vom feldfreien darstellen: sie bilden einen Tensor vom zweiten Range gegenüber Lorentz-Transformationen.

§ 1. Lösung der Näherungsgleichungen des Gravitationsfeldes durch retardierte Potentiale.

Wir gehen aus von den für ein beliebiges Koordinatensystem gültigen² Feldgleichungen

$$-\sum_{\alpha} \frac{\partial}{\partial x_{\alpha}} \left\{ \frac{\mu_{\alpha}}{x} \right\} + \sum_{\alpha} \frac{\partial}{\partial x_{\alpha}} \left\{ \frac{u_{\alpha}}{x} \right\} + \sum_{\alpha} \left\{ \frac{\mu_{\alpha}}{\delta} \right\} \left\{ \frac{\delta}{x} \right\} - \sum_{\alpha\beta} \left\{ \frac{u_{\alpha\beta}}{x} \right\} \left\{ \frac{x}{\delta} \right\} \quad (2)$$

$$= -\kappa \left(\mathcal{T}_{\alpha\beta} - \frac{1}{2} g_{\alpha\beta} T \right)$$

¹ Diese Sitzungsber. 1916, S. 685f.

² Von der Einführung des δ -Gliedes (vgl. diese Sitzungsber. 1917, S. 149) ist dabei Abstand genommen.

AMPLITUDE AND POWER OF A GRAVITATIONAL WAVE

If R , M , v are the characteristic size, mass, and speed of the emitting source, and r is our distance from the source, we have:

Amplitude:
$$h \sim \frac{r_{\text{Sch}}}{r} \frac{v^2}{c^2},$$

Where r_{Sch} is the Schwarzschild radius of the emitting source: $r_{\text{Sch}} = GM/c^2$

Luminosity:
$$\frac{dE}{dt} \sim \frac{G}{c^5} \left(\frac{M}{R} \right)^2 v^6 \sim L_0 \left(\frac{r_{\text{Sch}}}{R} \right)^2 \left(\frac{v}{c} \right)^6,$$

where

$$L_0 \equiv \frac{c^5}{G} = 3.6 \times 10^{59} \text{ erg s}^{-1}.$$

Simulation:

Manuela Campanlli

Carlos Lousto

Yosef Zlochower

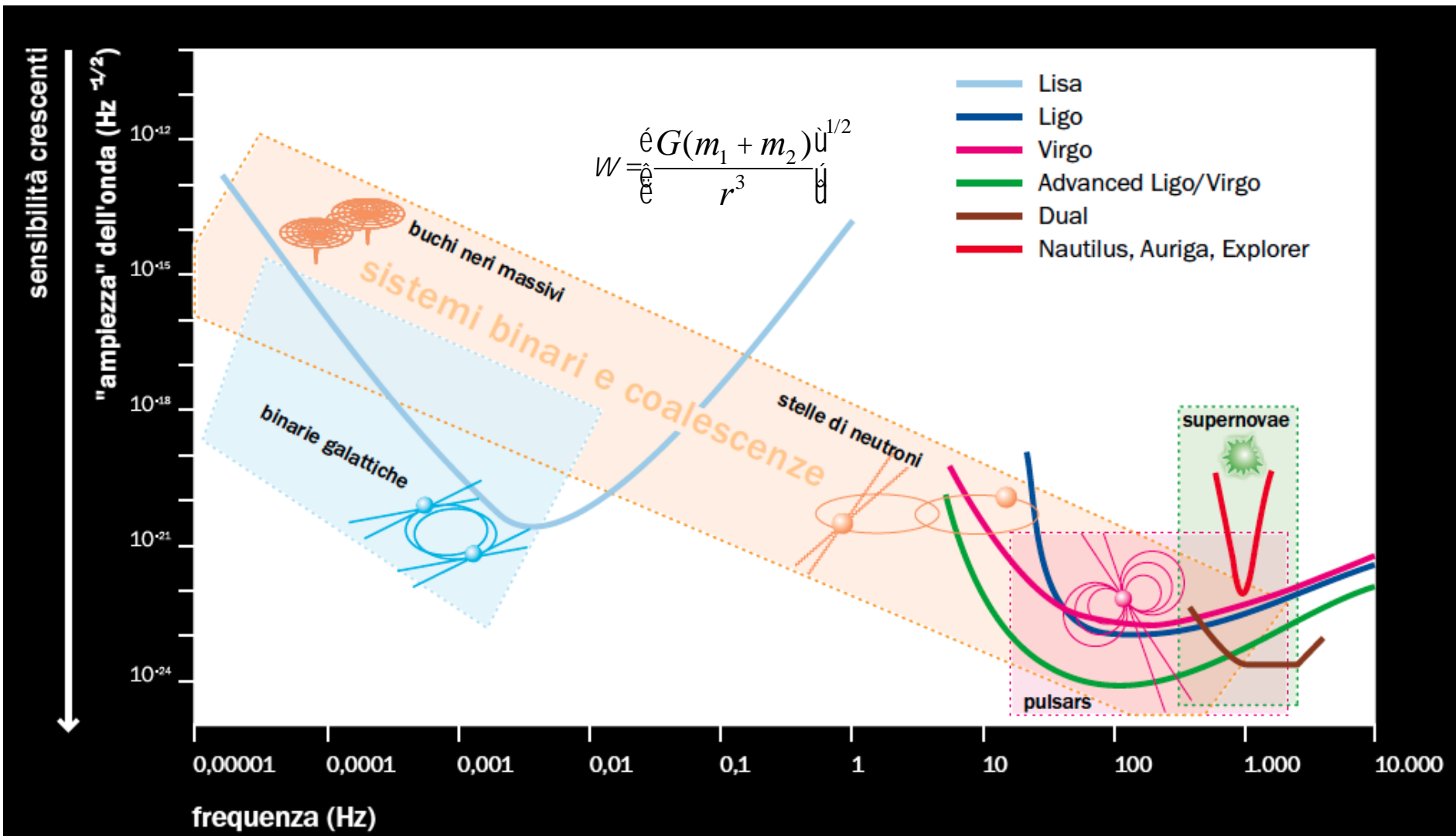
Visualization:

Hans-Peter Bischof

CCRG
RIT

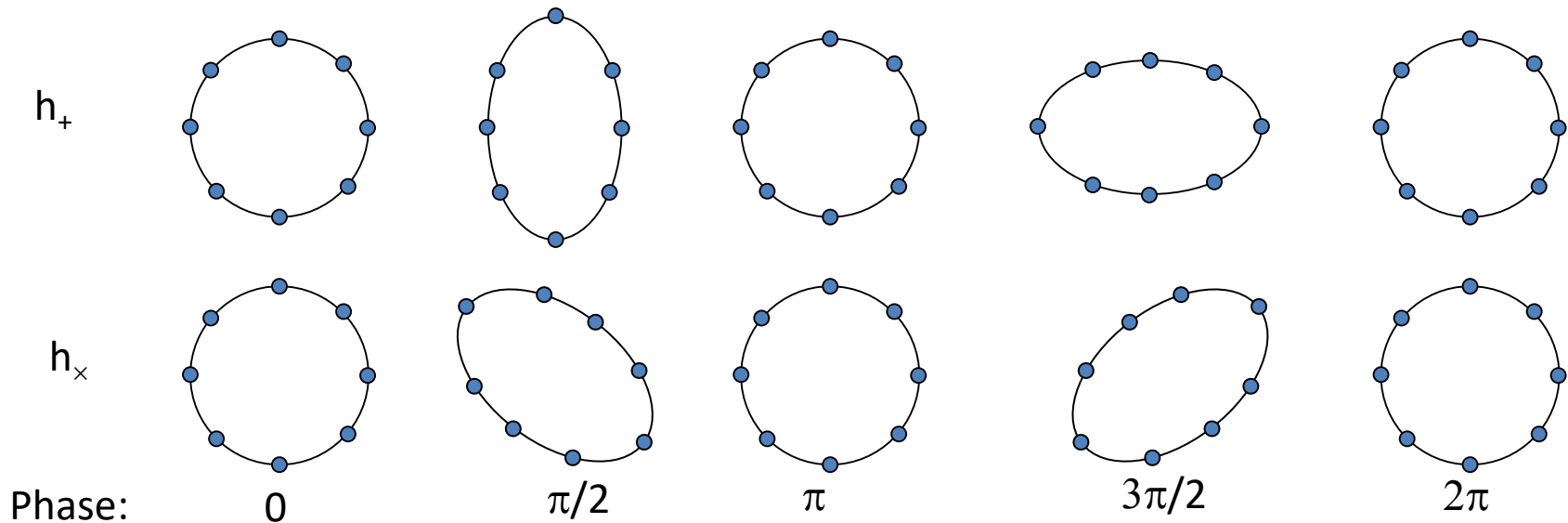
Copyright - CCRG - 2009

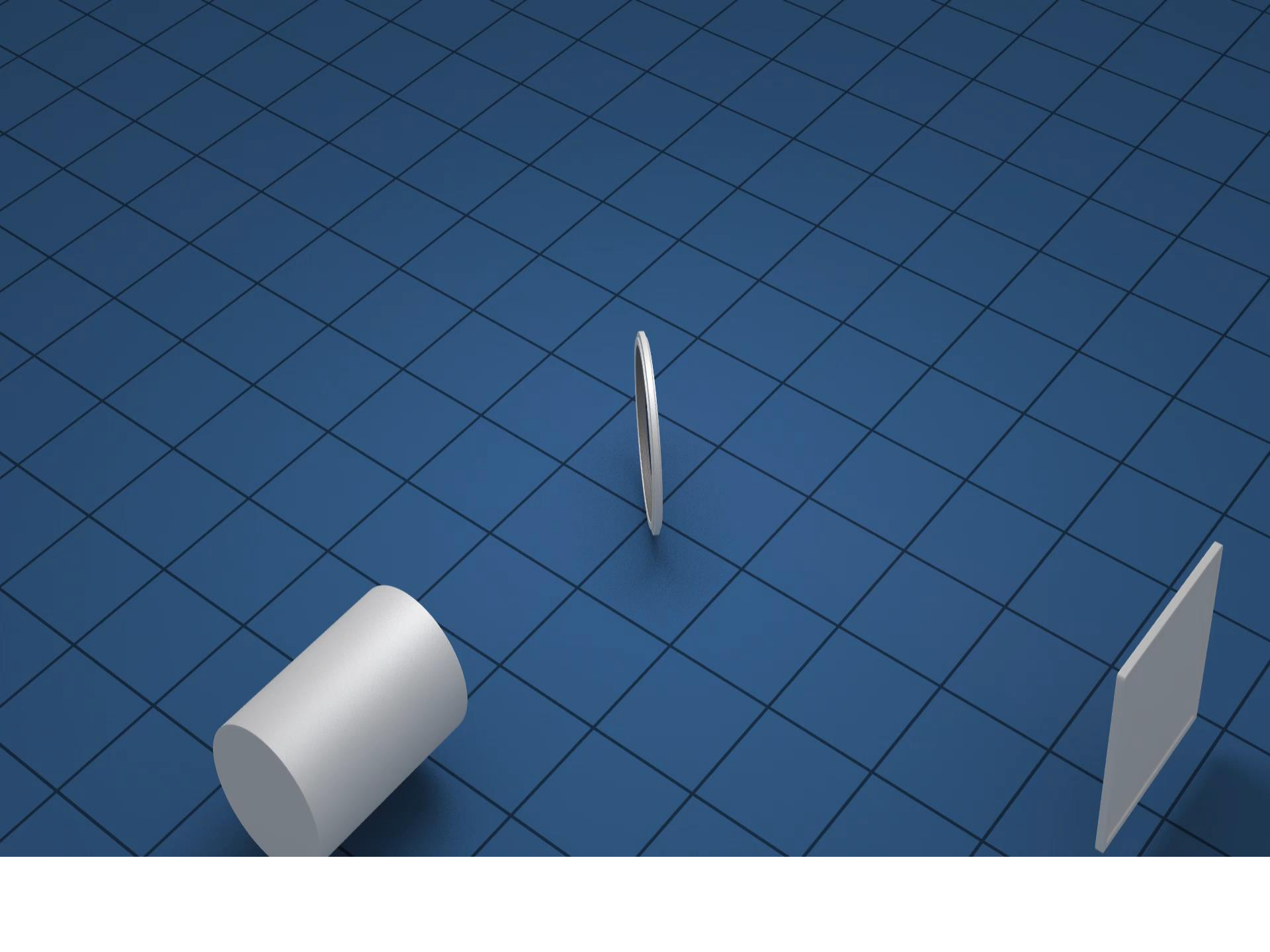
GW amplitude-frequency diagram of known cosmic sources



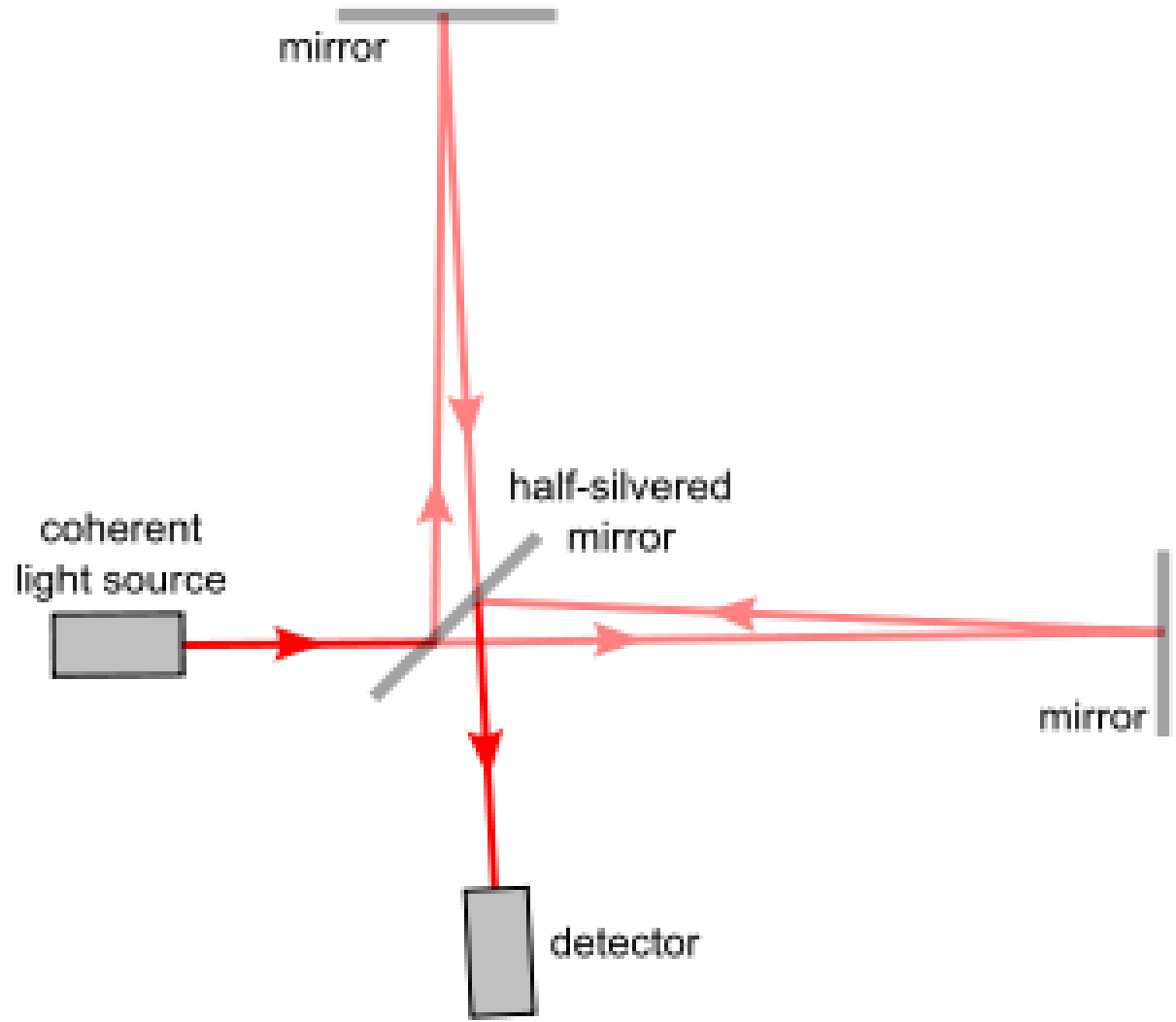
Polarization of GWs

GWs have 2 different linear polarizations. The 2 states are rotated 45 deg relative to one another. This contrasts with the 2 polarization states of an electromagnetic wave, which are 90 deg to each other. This pattern of polarization is due to the fact that gravity is represented by a second-rank symmetric **tensor** ($h_{\mu\nu}$), while electromagnetism is represented by the **vector** potential A^{mu}





Scheme of a laser interferometer



$$h = \frac{\Delta L}{L}$$



LIGO, Hanford, Washington State, USA



KAGRA (Japan)

**Fourth Observing Run (O4):
May 2023**

GRAVITATIONAL INTERFEROMETER NETWORK: LIGO, VIRGO, KAGRA



LIGO, Livingston, Louisiana, USA



**VIRGO
Cascina (Italy)**

Advanced LIGO
Hanford
2015



GEO600 (HF)
2011



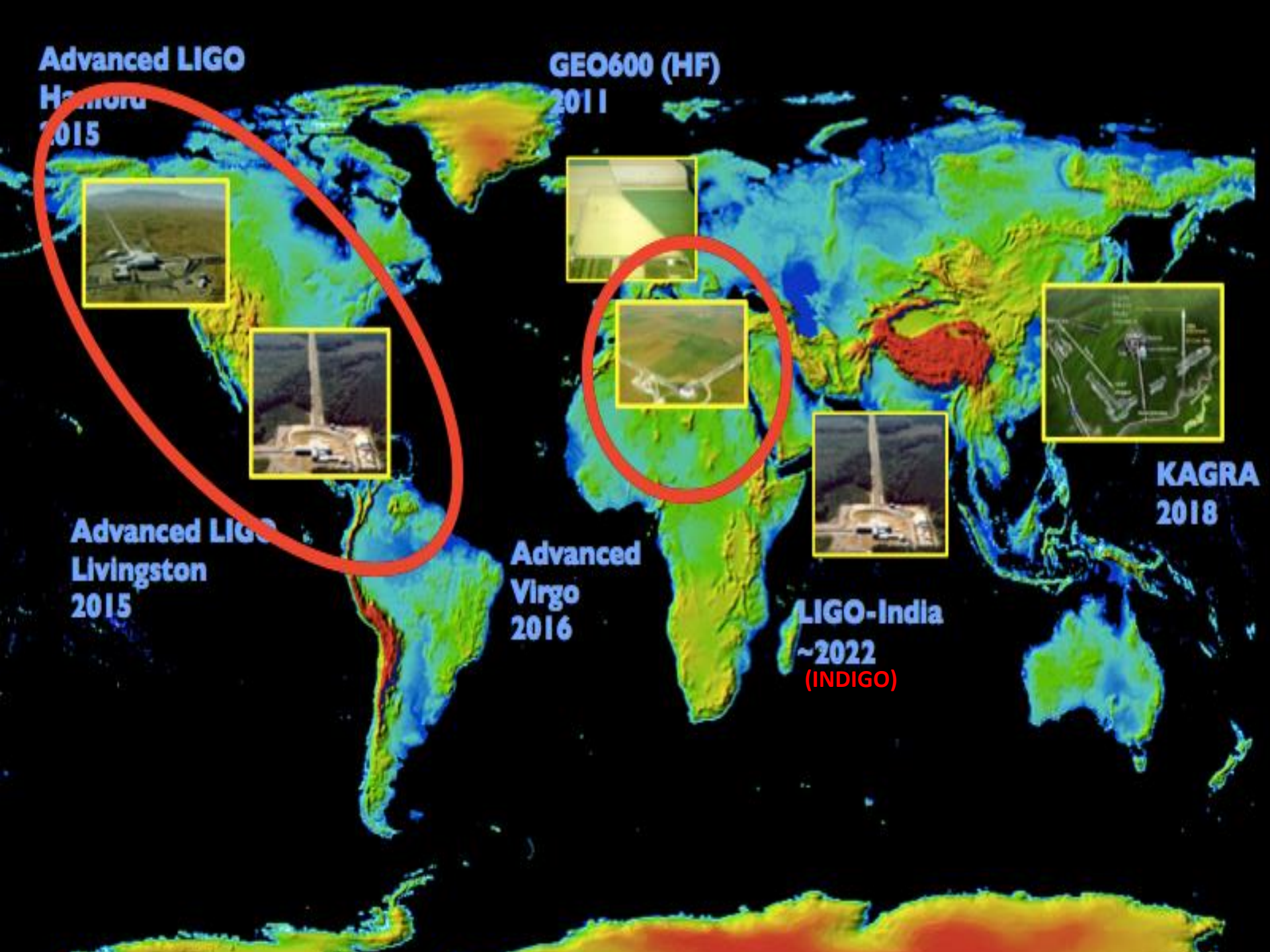
KAGRA
2018

Advanced LIGO
Livingston
2015

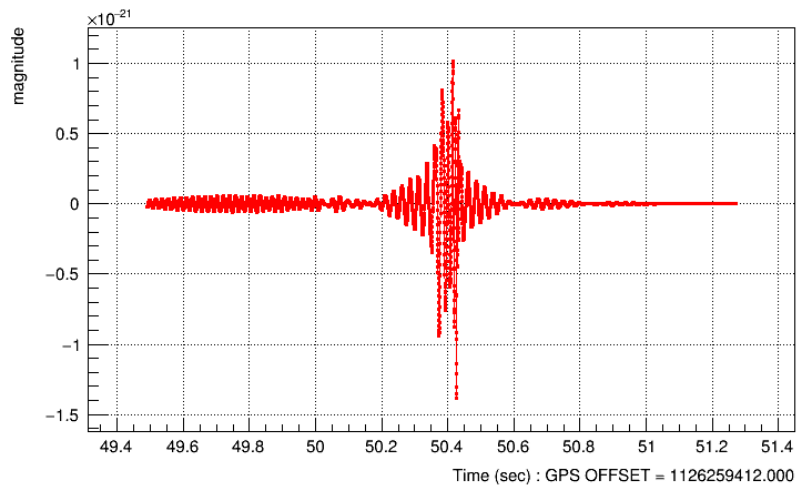
Advanced Virgo
2016



LIGO-India
~2022
(INDIGO)

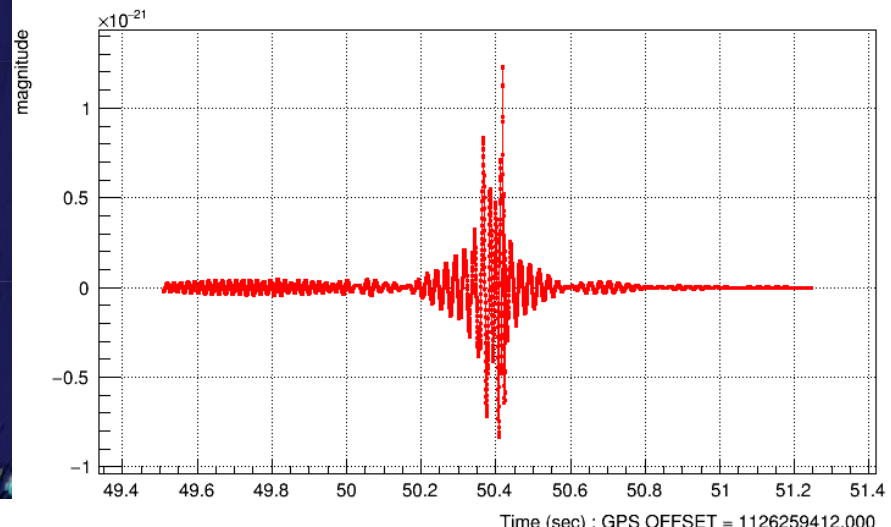
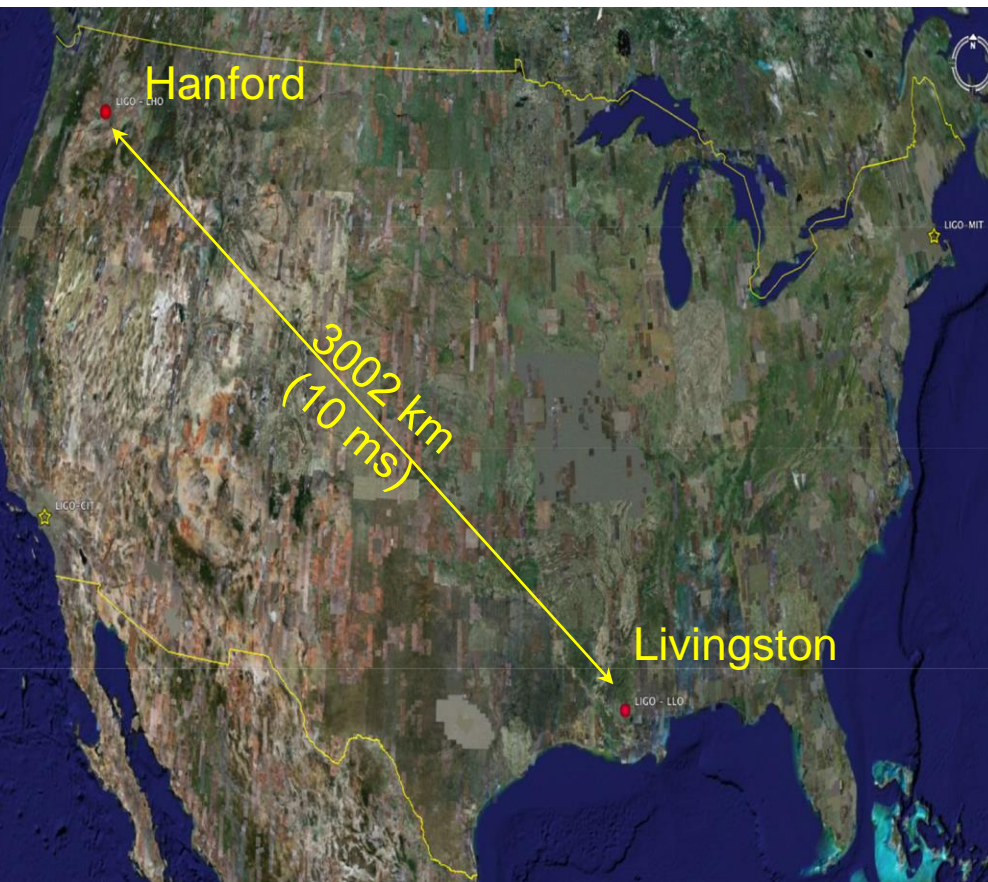


**September 14, 2015 at
11:50:45 in Central
European Time**

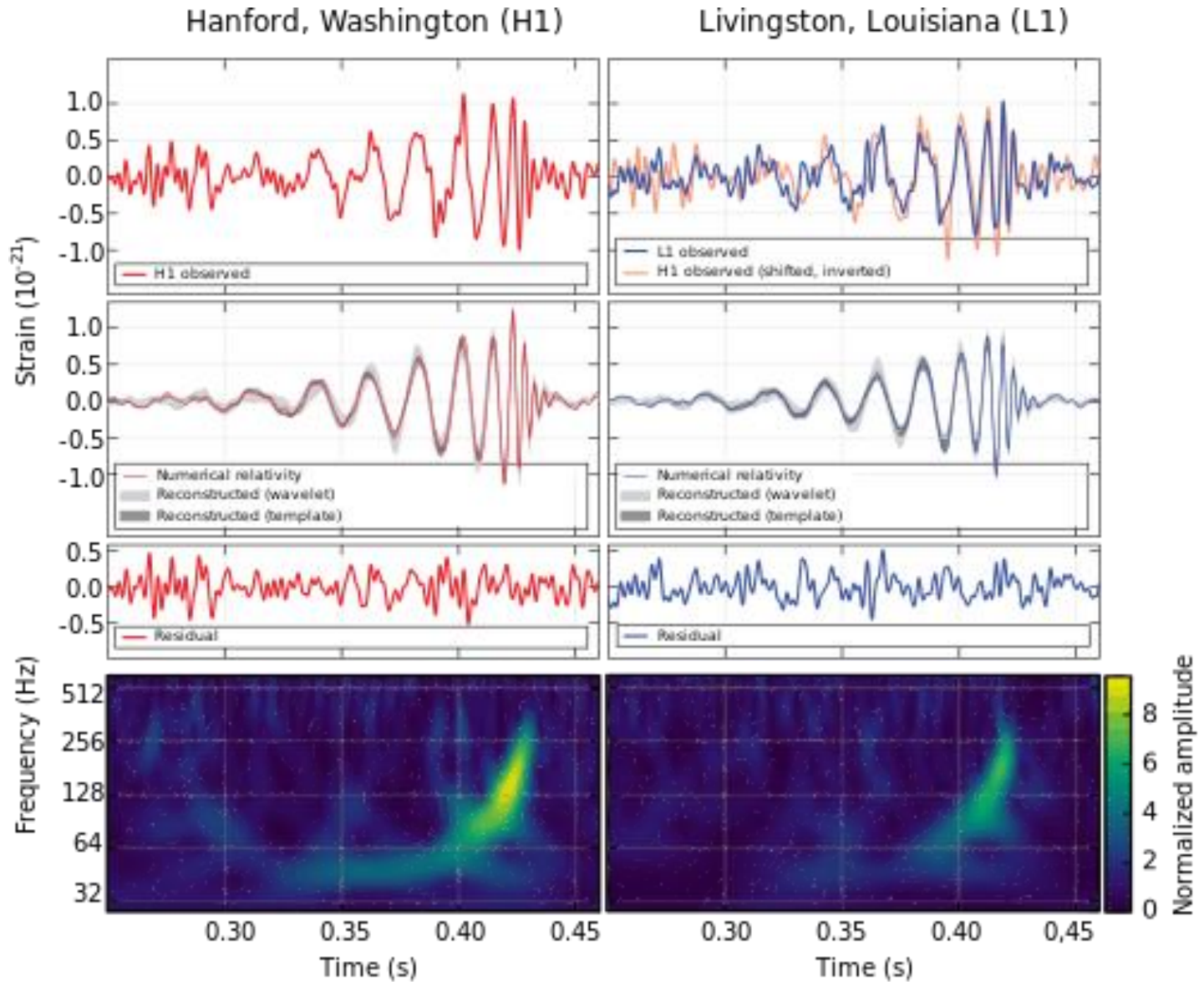


Alarm reported by the on-line algorithm for generic transient search

SNR = 24



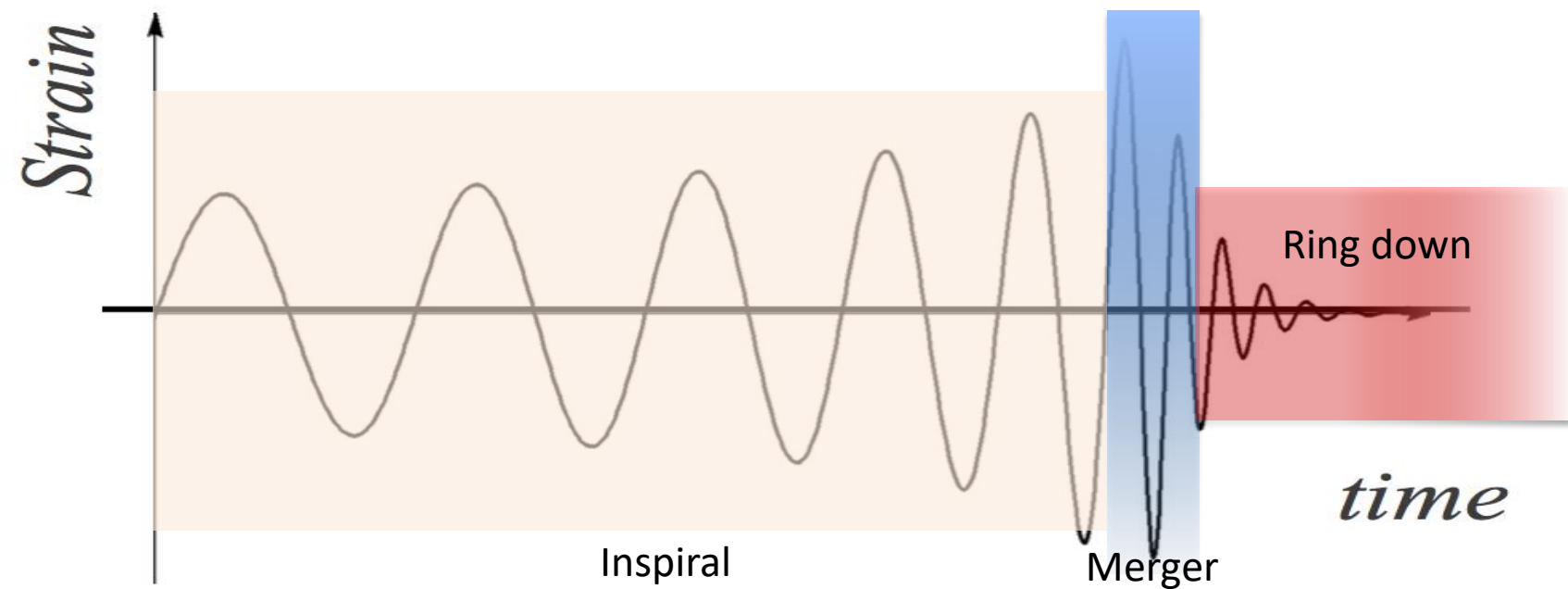
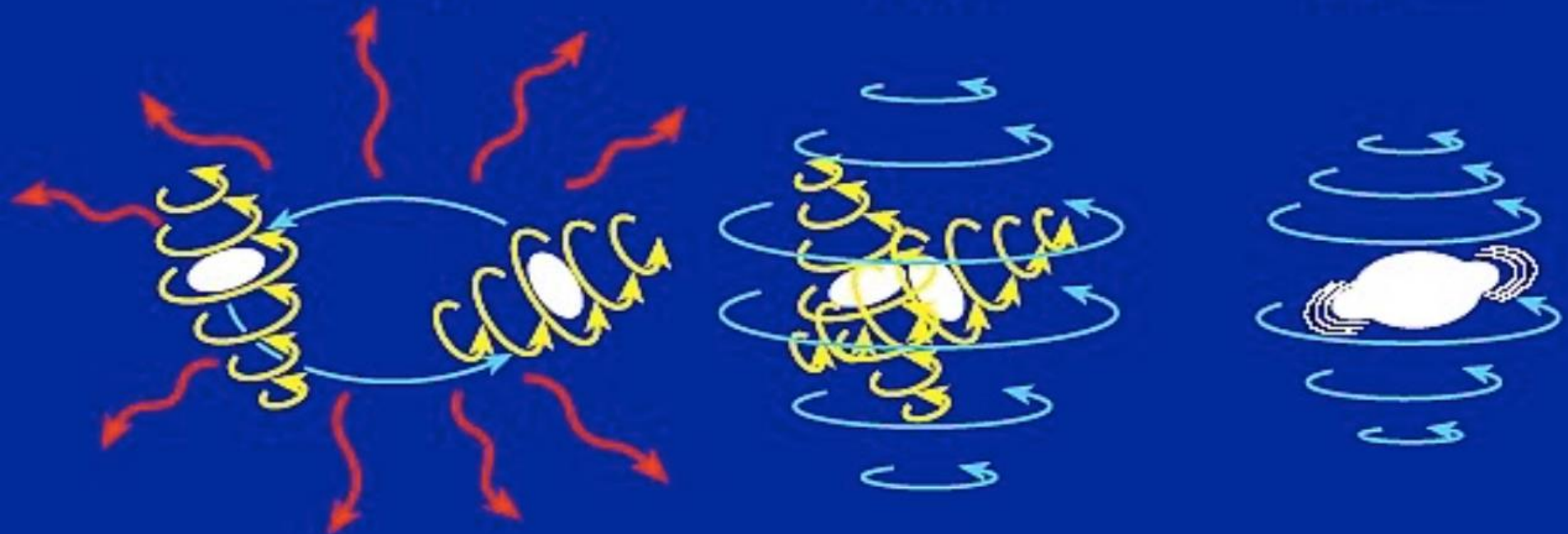
14 September 2015: First detection of gravitational waves



Inspiral

Merger

Ringdown





Observation of Gravitational Waves from a Binary Black Hole Merger

B. P. Abbott *et al.**

(LIGO Scientific Collaboration and Virgo Collaboration)

(Received 21 January 2016; published 11 February 2016)

On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed a transient gravitational-wave signal. The signal sweeps upwards in frequency from 35 to 250 Hz with a peak gravitational-wave strain of 1.0×10^{-21} . It matches the waveform predicted by general relativity for the inspiral and merger of a pair of black holes and the ringdown of the resulting single black hole. The signal was observed with a matched-filter signal-to-noise ratio of 24 and a false alarm rate estimated to be less than 1 event per 203 000 years, equivalent to a significance greater than 5.1σ . The source lies at a luminosity distance of 410_{-180}^{+160} Mpc corresponding to a redshift $z = 0.09_{-0.04}^{+0.03}$. In the source frame, the initial black hole masses are $36_{-4}^{+5} M_{\odot}$ and $29_{-4}^{+4} M_{\odot}$, and the final black hole mass is $62_{-4}^{+4} M_{\odot}$, with $3.0_{-0.5}^{+0.5} M_{\odot} c^2$ radiated in gravitational waves. All uncertainties define 90% credible intervals.

These observations demonstrate the existence of binary stellar-mass black hole systems. This is the first direct detection of gravitational waves and the first observation of a binary black hole merger.

DOI: 10.1103/PhysRevLett.116.061102

<http://link.aps.org/doi/10.1103/PhysRevLett.116.061102>

844 Sitzung der physikalisch-mathematischen Klasse vom 25. November 1915

Die Feldgleichungen der Gravitation.

Von A. EINSTEIN.

In zwei vor kurzem erschienenen Mitteilungen¹ habe ich gezeigt, wie man zu Feldgleichungen der Gravitation gelangen kann, die dem Postulat allgemeiner Relativität entsprechen, d. h. die in ihrer allgemeinen Fassung beliebigen Substitutionen der Raumkoordinaten gegenüber kovariant sind.

Die Entwicklungsgang war dabei folgender. Zunächst fand ich Gleichungen, welche die Newtonsche Theorie als Näherung enthalten und beliebigen Substitutionen von drei Determinanten 1. ordnung Kovariant waren. Hierauf fand ich, daß diese Gleichungen allgemein kovariant entsprechen, falls der Skalar des Energietensors der Materie verschwindet. Das Koordinatensystem war dann nach der einfachen Regel zu spezialisieren, daß $1 = g_{11}$ gesetzt wird, wodurch die Gleichungen der Theorie eine endliche Vereinfachung erfahren. Dabei zeigte aber, wie erwähnt, die Hypothese eingeführt werden, daß der Skalar des Energietensors der Materie verschwindet.

Neuzeitliche Studie ich nun, daß man ohne Hypothese über den Energietensor der Materie nachkommen kann, wenn man den Energietensor der Materie in etwas anderer Weise in die Feldgleichungen einsetzt, als dies in meinen beiden früheren Mitteilungen geschehen ist. Die Feldgleichungen für das Vakuum, auf welche ich die Erklärung der Perihelienbewegung des Merkur gegründet habe, bleiben von dieser Modifikation unberührt. Ich gebe hier zunächst die ganze Beschreibung, damit der Leser nicht gezwungen ist, die früheren Mitteilungen unzusammenhängend kennenzulernen.

Aus der bekannten RIEMANNSCHEN Kovariante vierter Ranges läßt sich folgende Kovariante vierter Ranges ableiten:

$$G_{\alpha\beta} = R_{\alpha\beta} + S_{\alpha\beta} \quad (1)$$

$$R_{\alpha\beta} = -\frac{1}{2} \sum_{\gamma} \frac{\partial^2 \Gamma^{\gamma}_{\alpha\beta}}{\partial x^{\gamma} \partial x^{\gamma}} + \sum_{\gamma} \frac{\partial \Gamma^{\gamma}_{\alpha\beta}}{\partial x^{\gamma}} \Gamma^{\gamma} \quad (2)$$

$$S_{\alpha\beta} = \sum_{\gamma} \frac{\partial^2 \Gamma^{\gamma}_{\alpha\beta}}{\partial x^{\gamma} \partial x^{\gamma}} - \sum_{\gamma} \frac{\partial \Gamma^{\gamma}_{\alpha\beta}}{\partial x^{\gamma}} \Gamma^{\gamma} \quad (3)$$

¹ Sitzungsber. XLIV, S. 178 und XLV, S. 179, 1915.

Über Gravitationswellen.

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Die wichtige Frage, wie die Ausbreitung der Gravitationsfelder erfolgt, ist schon vor anderthalb Jahren in einer Akademiesitzung von mir behandelt worden¹. In der meine damalige Darstellung des Gegenstandes nicht genügend durchsichtig und außerdem durch einen technischen Rechenfehler verunstaltet ist, muß ich hier nochmals auf die Angelegenheit zurückkommen.

Wie damals hervorgehoben ist, muß auch hier auf den Fall, daß das betrachtete zentralische System sich von einem „quasiisolierten“ sehr wenig unterscheiden. Um für alle Indizes

$$g_{\alpha\beta} = -\delta_{\alpha\beta} + \gamma_{\alpha\beta} \quad (1)$$

setzen zu können, wählen wir, wie es in der speziellen Relativitätstheorie üblich ist, die Zeitachse x_4 vom Ursprung, indem wir

$$x_4 = ct$$

setzen, wobei t die „Lichtzeit“ bedeutet. In (1) ist $\delta_{\alpha\beta} = 1$ bzw. $\delta_{\alpha\beta} = 0$, je nachdem $\alpha = \beta$ oder $\alpha \neq \beta$ ist. Die $\gamma_{\alpha\beta}$ sind gegen ϵ kleine Größen, welche die Abweichung des Kontinuum vom Feldfreizustand bilden einen Tensor vom zweiten Range gegenüber Lorentz-Transformationen.

§ 1. Lösung der Näherungsgleichungen des Gravitationsfeldes durch retardierte Potentiale.

Wir gehen aus von den für ein beliebiges Koordinatensystem gültigen Feldgleichungen

$$-\sum_{\alpha} \frac{\partial^2 \Gamma^{\alpha}_{\alpha\beta}}{\partial x^{\alpha} \partial x^{\alpha}} + \sum_{\alpha} \frac{\partial \Gamma^{\alpha}_{\alpha\beta}}{\partial x^{\alpha}} + \sum_{\alpha} \frac{\partial \Gamma^{\alpha}_{\beta\alpha}}{\partial x^{\alpha}} - \sum_{\alpha} \frac{\partial^2 \Gamma^{\alpha}_{\beta\alpha}}{\partial x^{\alpha} \partial x^{\alpha}} = -\kappa \left(T_{\alpha\beta} - \frac{1}{2} g_{\alpha\beta} T \right) \quad (2)$$

¹ Diese Sitzungsber. 1915, S. 688f.
² Von der Erklärung des „3-Gliedes“ (vgl. diese Sitzungsber. 1917, S. 141) ist ab hier Abstand genommen.

GW150914: FACTSHEET

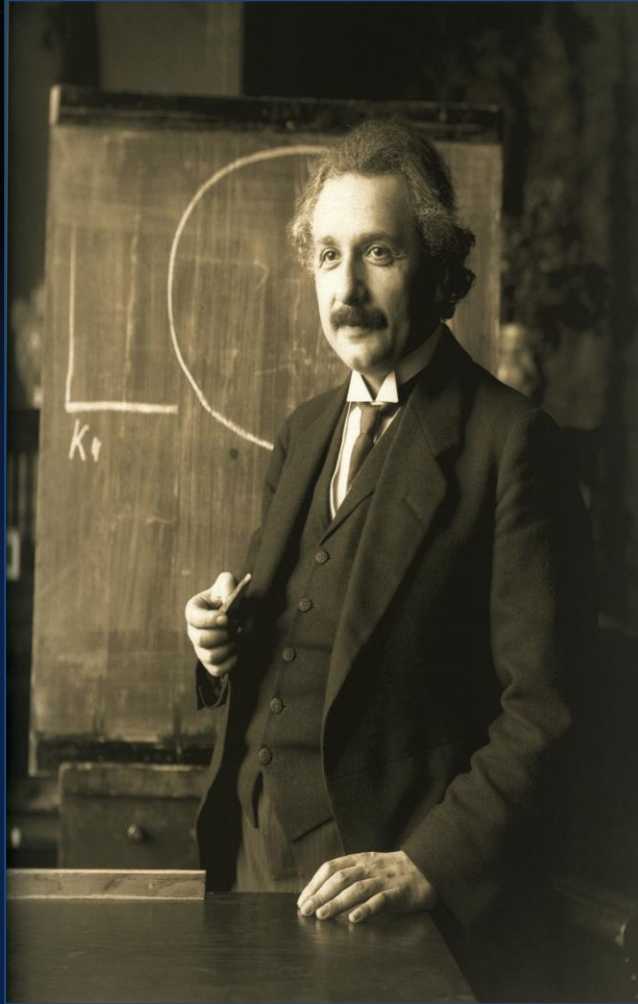
BACKGROUND IMAGES: TIME-FREQUENCY TRACE (TOP) AND TIME-SERIES (BOTTOM) IN THE TWO LIGO DETECTORS; SIMULATION OF BLACK HOLE HORIZONS (MIDDLE-TOP), BEST FIT WAVEFORM (MIDDLE-BOTTOM)

first direct detection of gravitational waves (GW) and first direct observation of a black hole binary

observed by	LIGO L1, H1	duration from 30 Hz	~ 200 ms
source type	black hole (BH) binary	# cycles from 30 Hz	~10
date	14 Sept 2015	peak GW strain	1×10^{-21}
time	09:50:45 UTC	peak displacement of interferometers arms	± 0.002 fm
likely distance	0.75 to 1.9 Gly 230 to 570 Mpc	frequency/wavelength at peak GW strain	150 Hz, 2000 km
redshift	0.054 to 0.136	peak speed of BHs	~ 0.6 c
signal-to-noise ratio	24	peak GW luminosity	3.6×10^{56} erg s ⁻¹
false alarm prob.	< 1 in 5 million	radiated GW energy	2.5-3.5 M _⊙
false alarm rate	< 1 in 200,000 yr	remnant ringdown freq.	~ 250 Hz
Source Masses	M _⊙	remnant damping time	~ 4 ms
total mass	60 to 70	remnant size, area	180 km, 3.5×10^5 km ²
primary BH	32 to 41	consistent with general relativity?	passes all tests performed
secondary BH	25 to 33	graviton mass bound	$< 1.2 \times 10^{-22}$ eV
remnant BH	58 to 67	coalescence rate of binary black holes	2 to 400 Gpc ⁻³ yr ⁻¹
mass ratio	0.6 to 1	online trigger latency	~ 3 min
primary BH spin	< 0.7	# offline analysis pipelines	5
secondary BH spin	< 0.9	CPU hours consumed	~ 50 million (=20,000 PCs run for 100 days)
remnant BH spin	0.57 to 0.72	papers on Feb 11, 2016	13
signal arrival time delay	arrived in L1 7 ms before H1	# researchers	~1000, 80 institutions in 15 countries
likely sky position	Southern Hemisphere		
likely orientation resolved to	face-on/off ~600 sq. deg.		

Detector noise introduces errors in measurement. Parameter ranges correspond to 90% credible bounds.
 Acronyms: L1=LIGO Livingston, H1=LIGO Hanford; Gly=giga lightyear= 9.46×10^{12} km; Mpc=mega parsec=3.2 million lightyear, Gpc= 10^3 Mpc, fm=femtometer= 10^{-15} m, M_⊙=1 solar mass= 2×10^{30} kg

Gravitational Waves



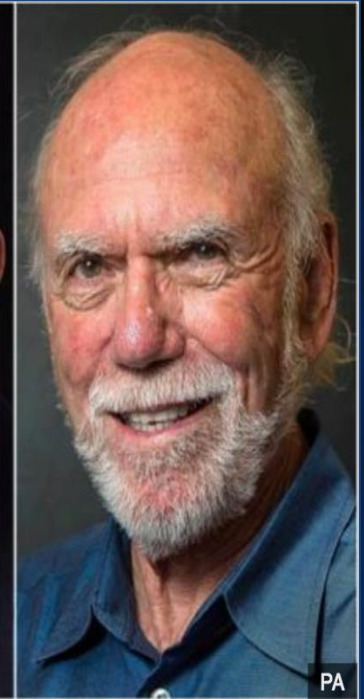
A. Einstein 1916



R. Weiss



K. Thorne



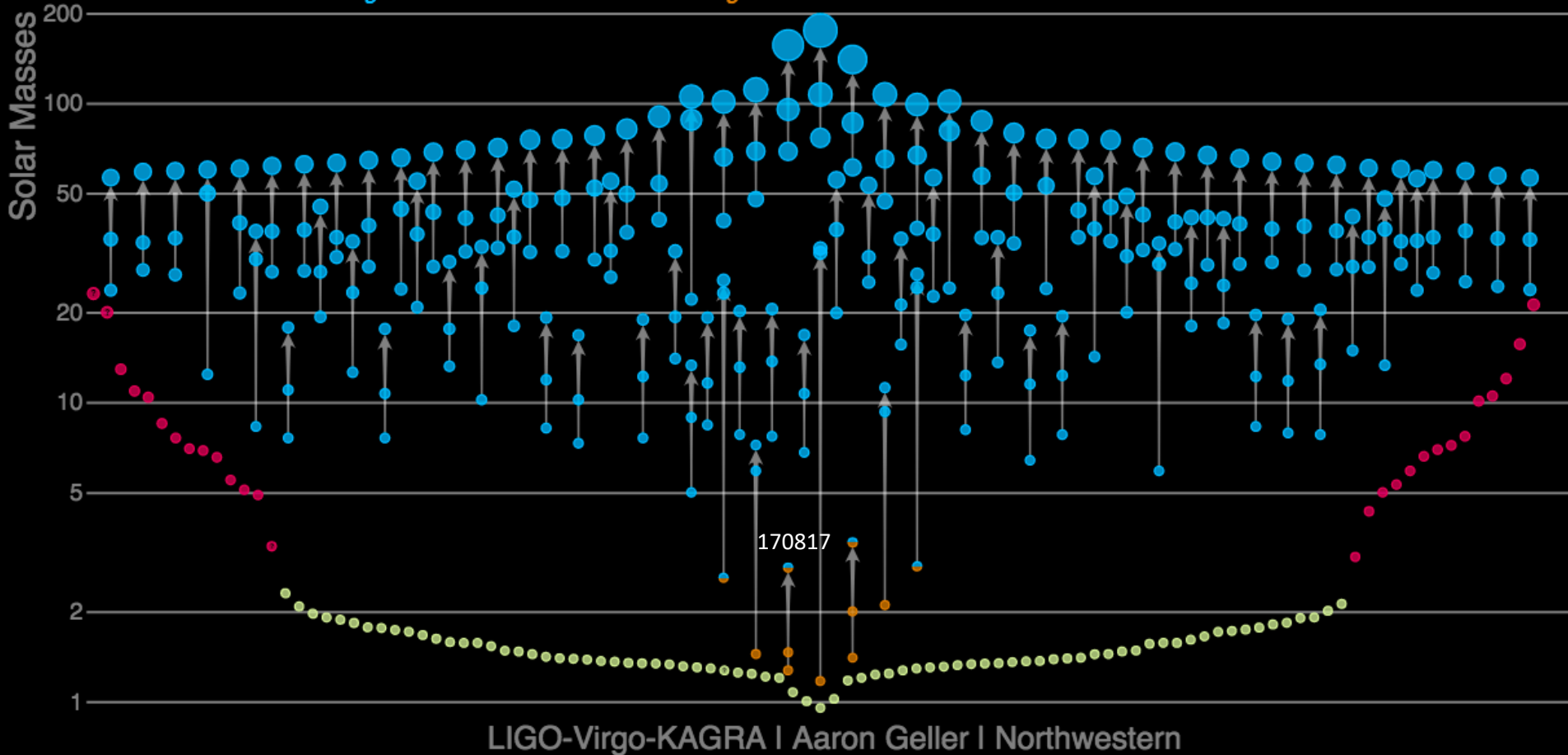
B. Barish

Nobel Prize for Physics 2017

Masses in the Stellar Graveyard



LIGO-Virgo-KAGRA Black Holes LIGO-Virgo-KAGRA Neutron Stars EM Black Holes EM Neutron Stars



<https://ligo.northwestern.edu/media/mass-plot/index.html>

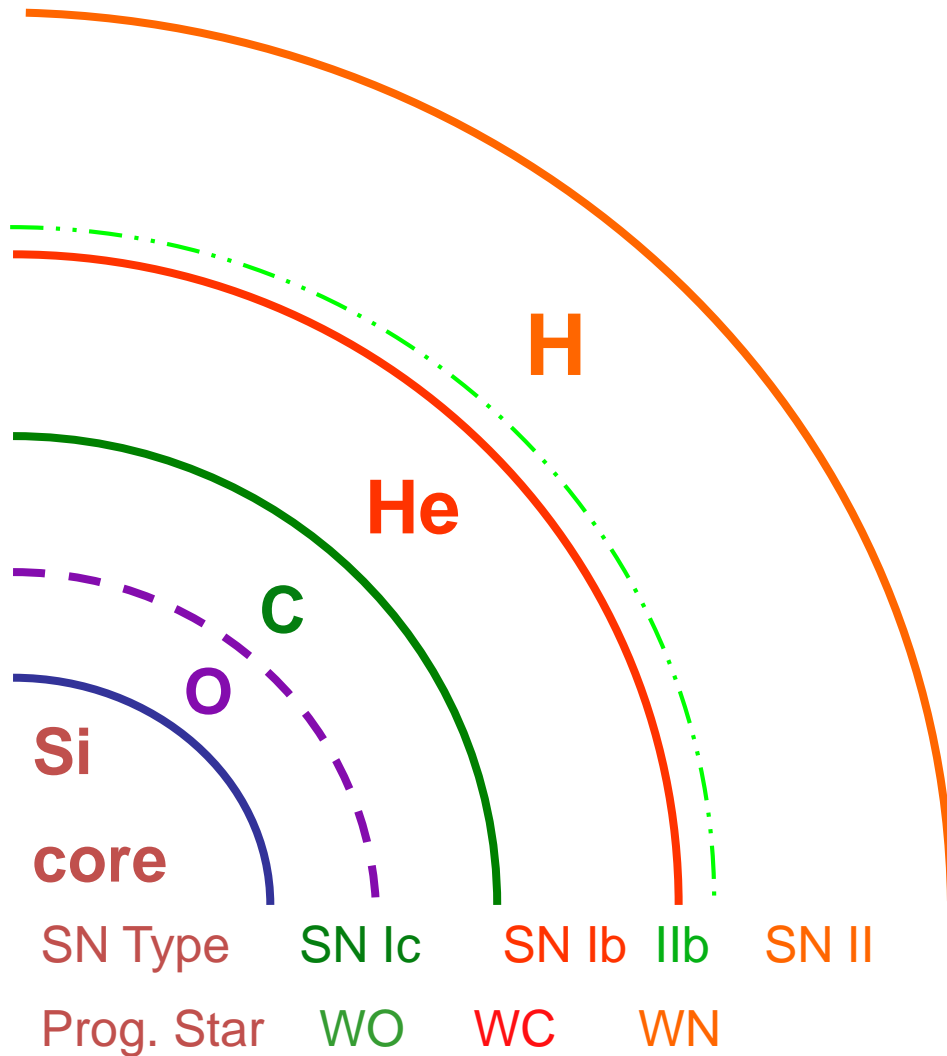
While black hole mergers are expected to be electromagnetically “silent”, neutron star mergers (or mergers of a NS and a BH) are expected to be accompanied by an electromagnetic signal. The mergers of systems containing at least a NS present thus a large potential for:

- 1) determining the NS equation of state
- 2) mapping the history of heavy elements formation
- 3) Studying the progenitor population of short GRBs

Mass of star determines its fate

Mass of Progenitor star	Remnant
$<3M_{\odot}$	He White Dwarf
3-6 M_{\odot}	C-O White Dwarf
6-8 M_{\odot}	O-Ne-Mg White Dwarf
8-23 M_{\odot}	Neutron Star
23--~140 M_{\odot}	Black Hole
$>\sim 140 M_{\odot}$	None?

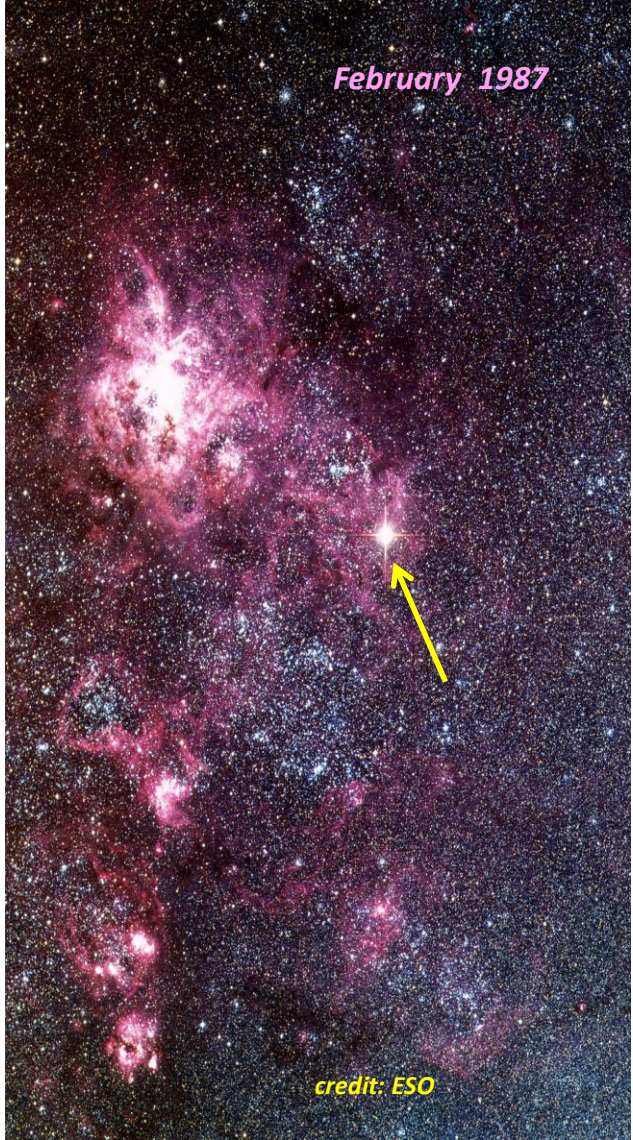
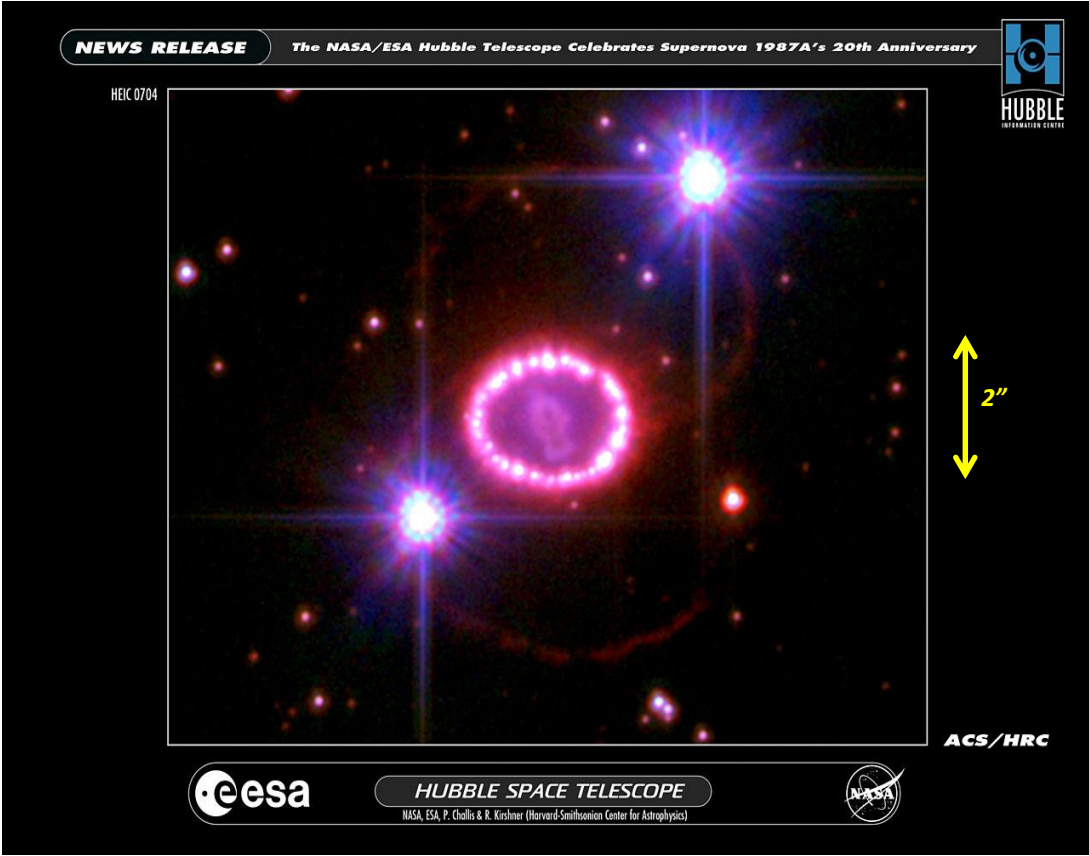
Massive Stars ($>8M_{\odot}$)



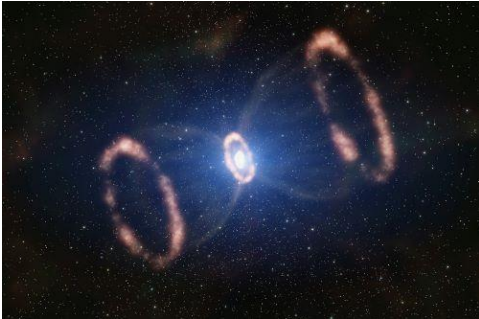
- Si burning \rightarrow Fe core
- **Core collapse**
- Compact object (**NS/BH**)
- ν emission
- KE deposited ($\sim 10^{51}$ erg)
- Nucleosynthesis
 \rightarrow ^{56}Ni ($\sim 0.1M_{\odot}$)
- envelope ejection
- SN type depends on degree of stripping
 - Role of binarity?

SN1987A: core-collapse supernova in the LMC (50 kpc)

First example of multi-messenger astronomy (barring the Sun...)



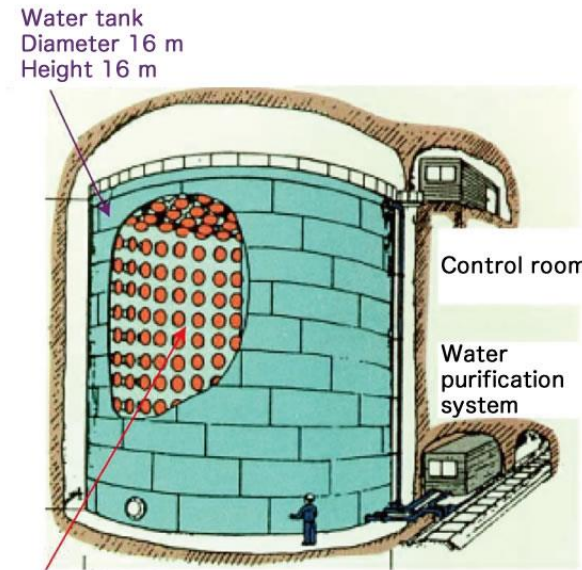
artist concept of SN1987A



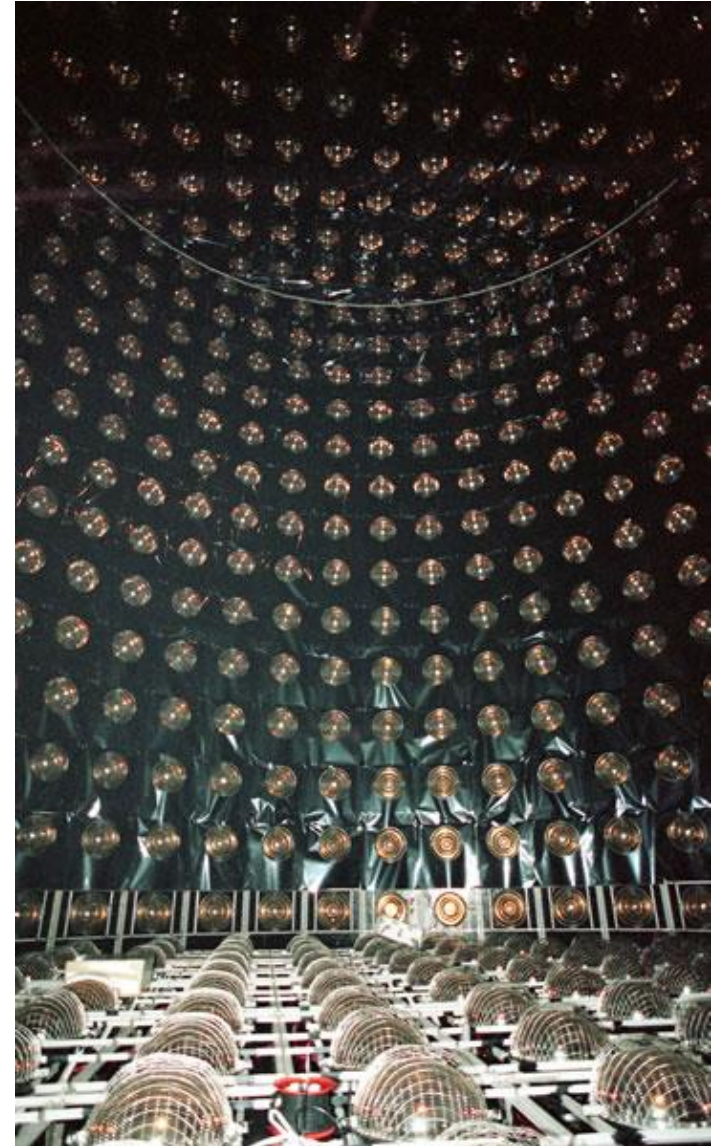


SN1987A: MeV neutrino detection Proves core-collapse and neutron star formation

The neutrino detectors IMB (USA), Kamiokande (Japan) and Baksan (Russia) detected a handful of neutrinos each, a few hours before light detection from SN1987A (I. Shelton 1987, IAU Circ. 4316)

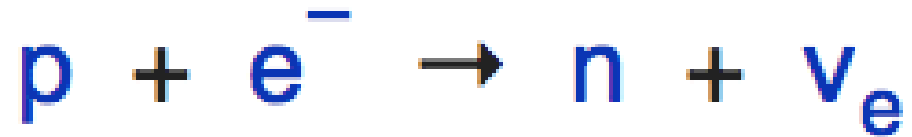


1000 20 inch
Photomultiplier Tubes



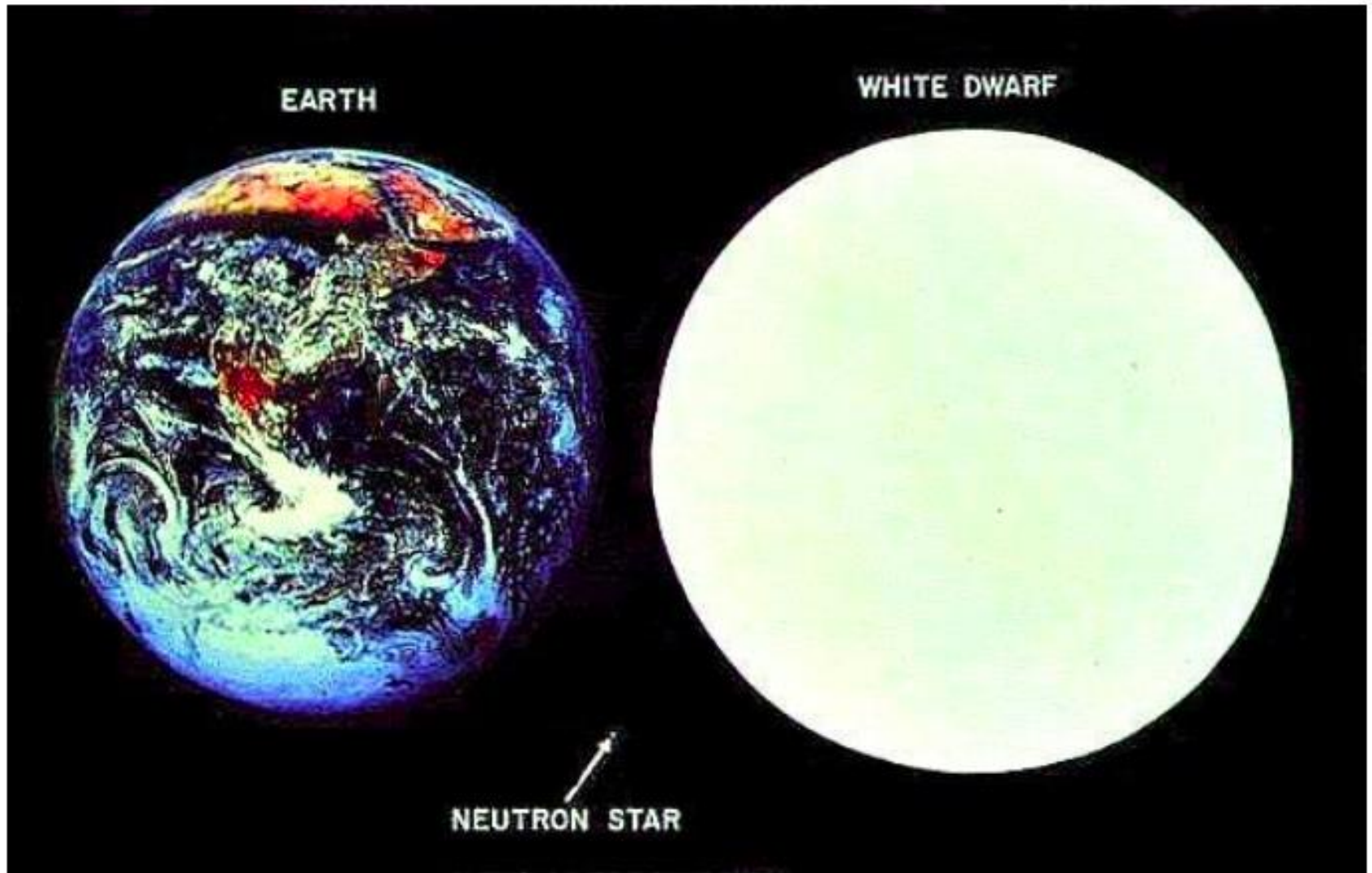
Core-collapse supernova neutrinos

Electron capture (or inverse beta decay):

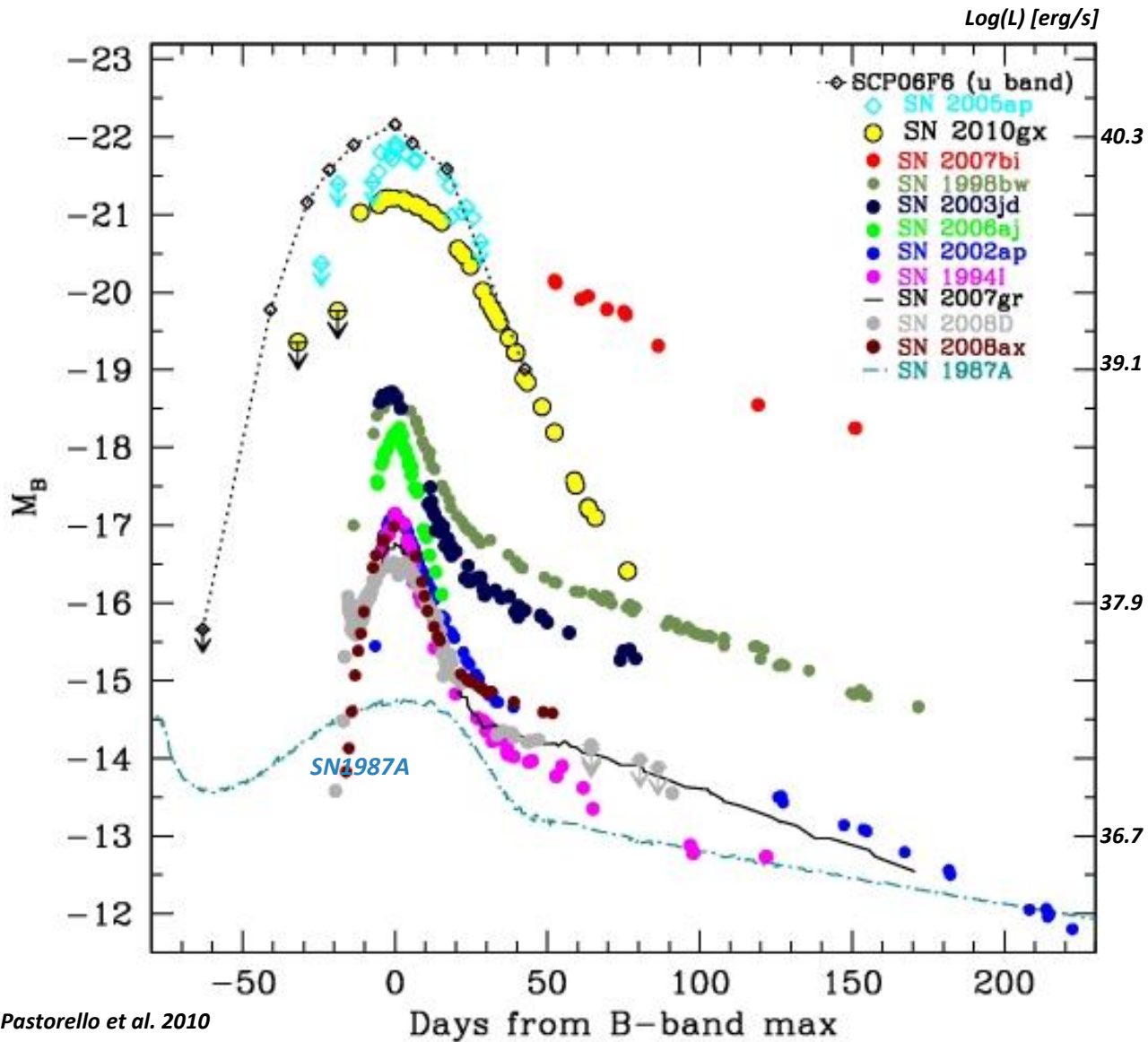


This process is responsible for neutronization of massive stellar cores undergoing gravitational collapse (neutron star remnant formation)

neutron star

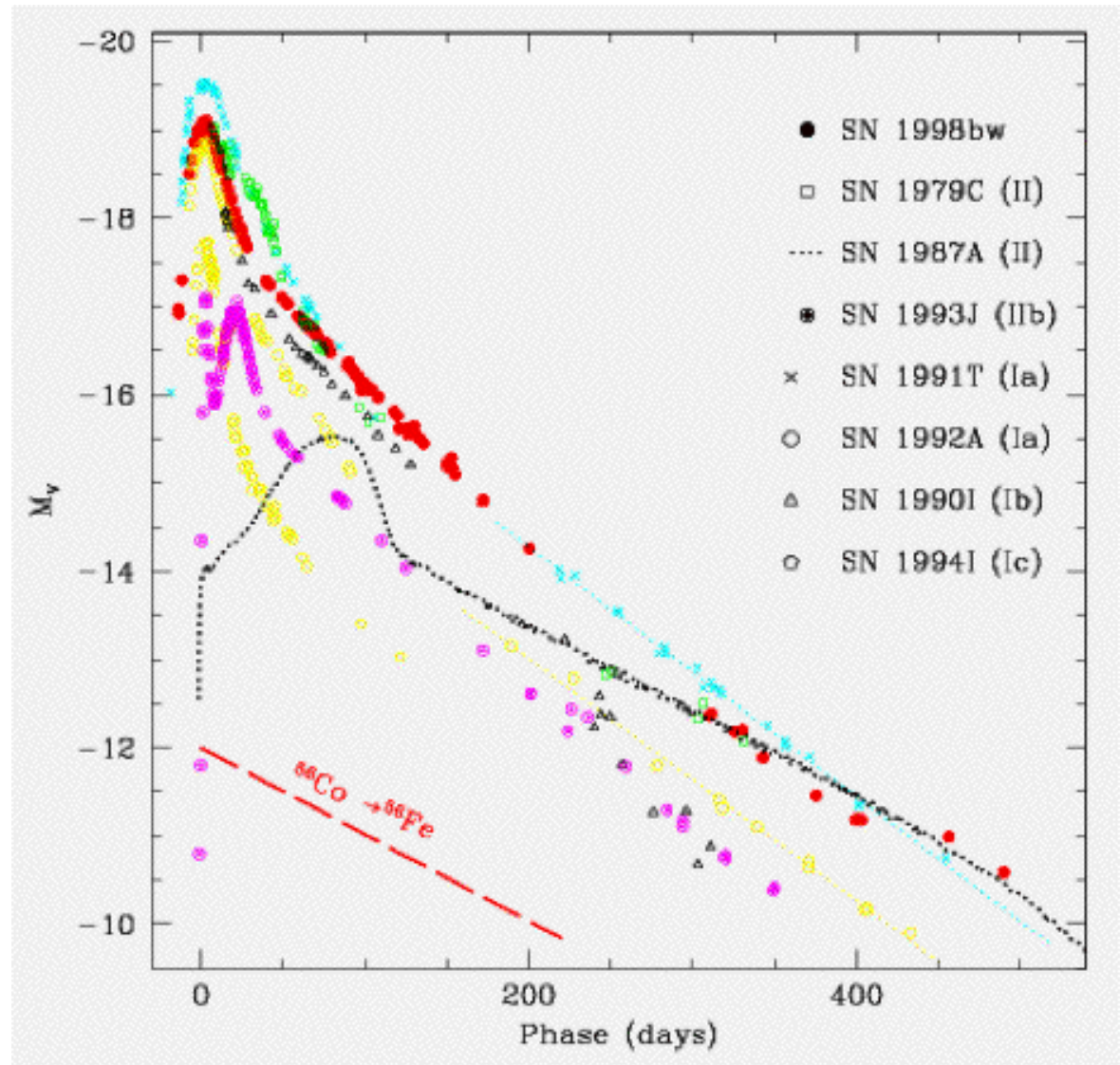


Light curves of core-collapse supernovae

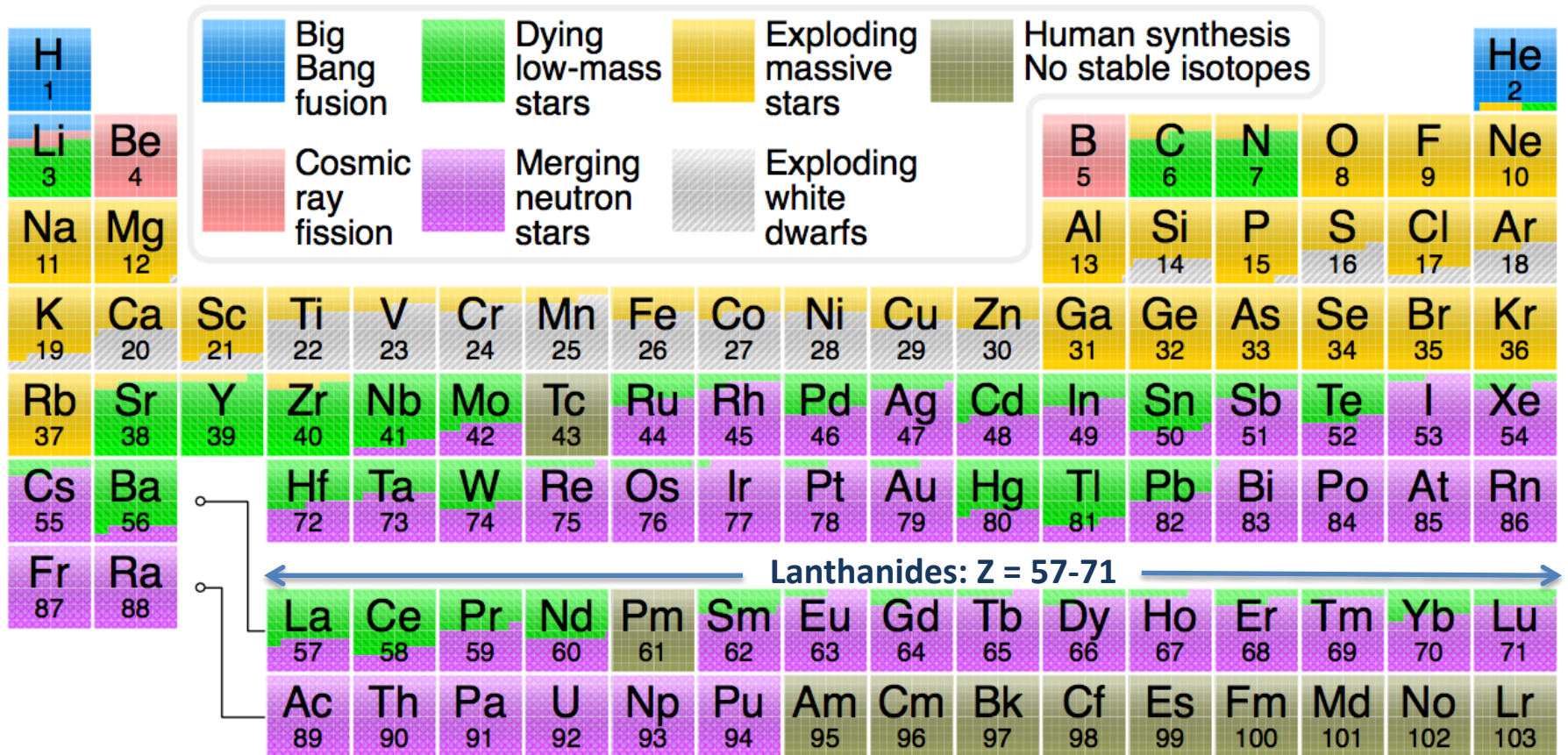


SN Light Curves are diverse

Powered
by
decay of
 ^{56}Ni
 $\rightarrow ^{56}\text{Co}$
 $\rightarrow ^{56}\text{Fe}$



Periodic table of elements



<https://en.wikipedia.org/wiki/R-process>

Electron capture ($p + e^- \rightarrow n + \nu_e$)

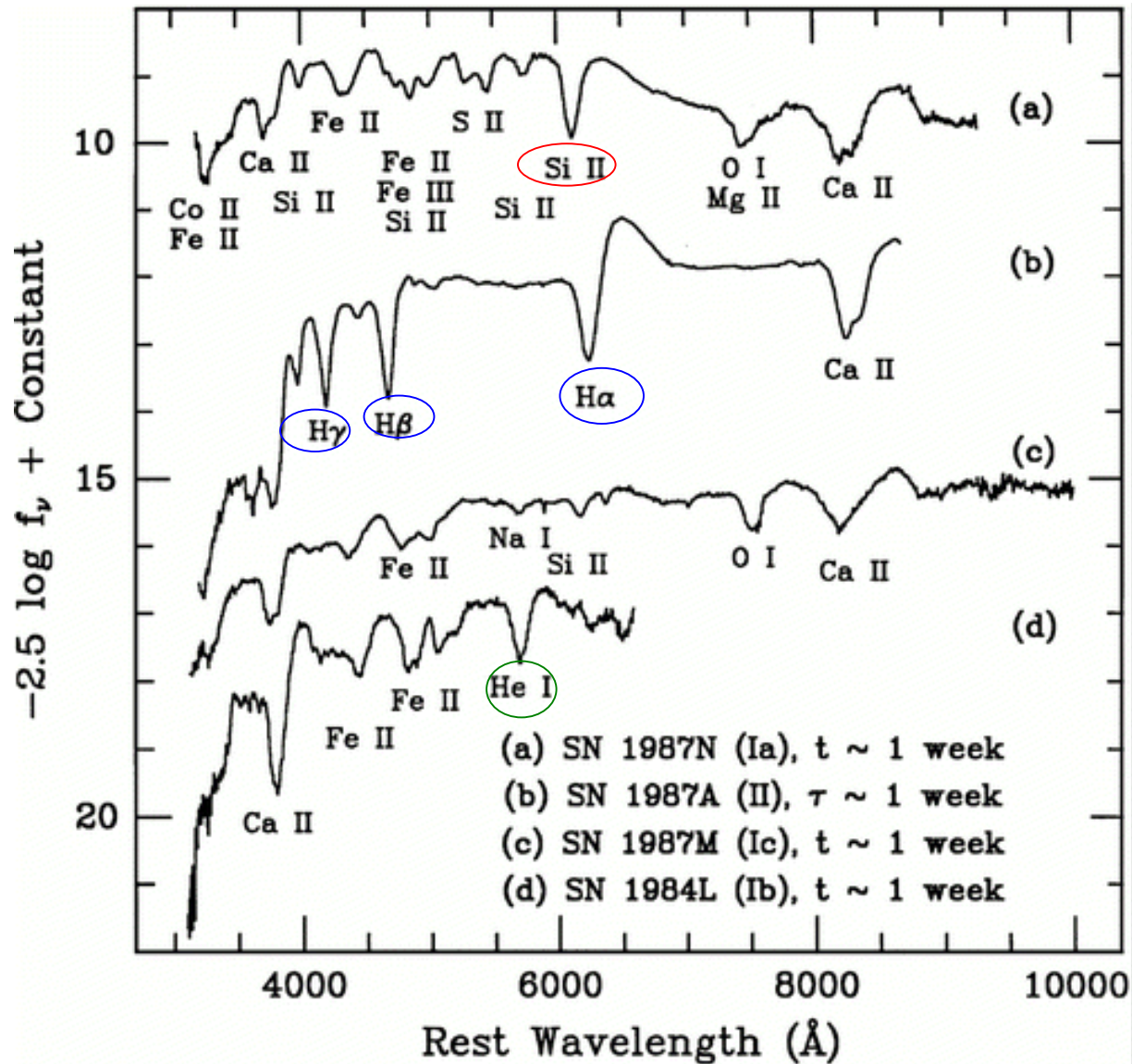
and beta+ decay ($p \rightarrow n + e^+ + \nu_e$)

Are also responsible for supernova light emission, dominated by the decay chain:



The emitted neutrino flux is normally way too weak to be detected

SN spectra near maximum light: diversity suggests different origin



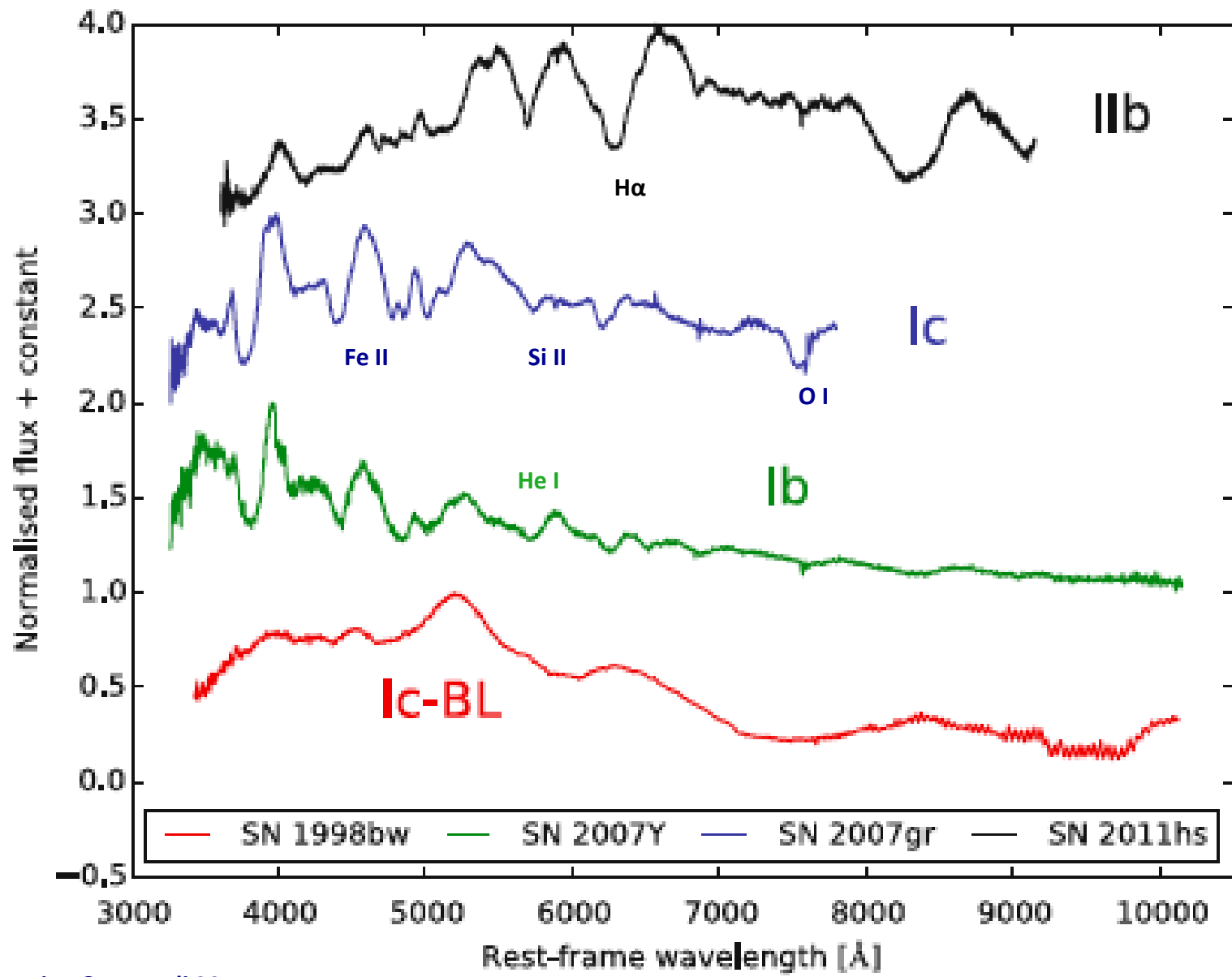
SN Ia

SN II

SN Ic

SN Ib

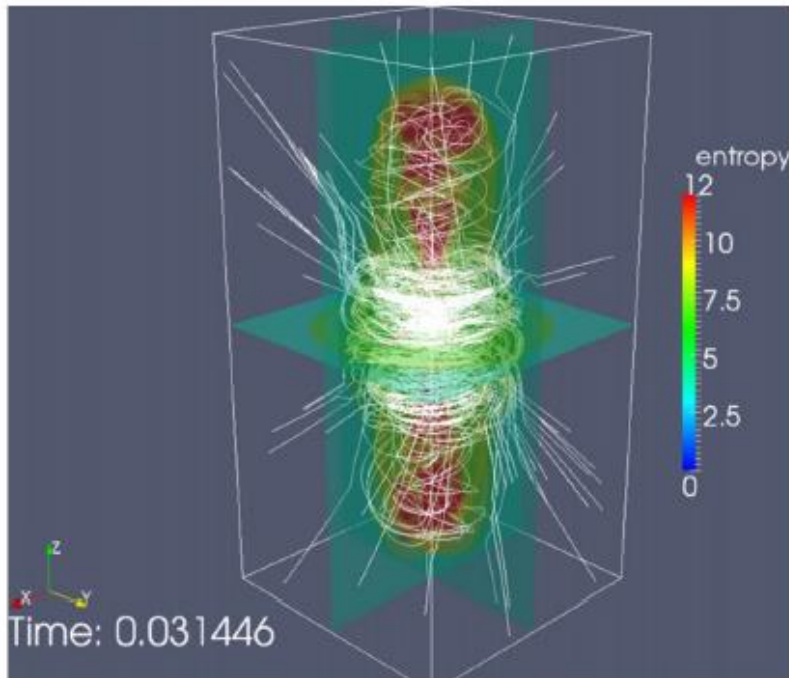
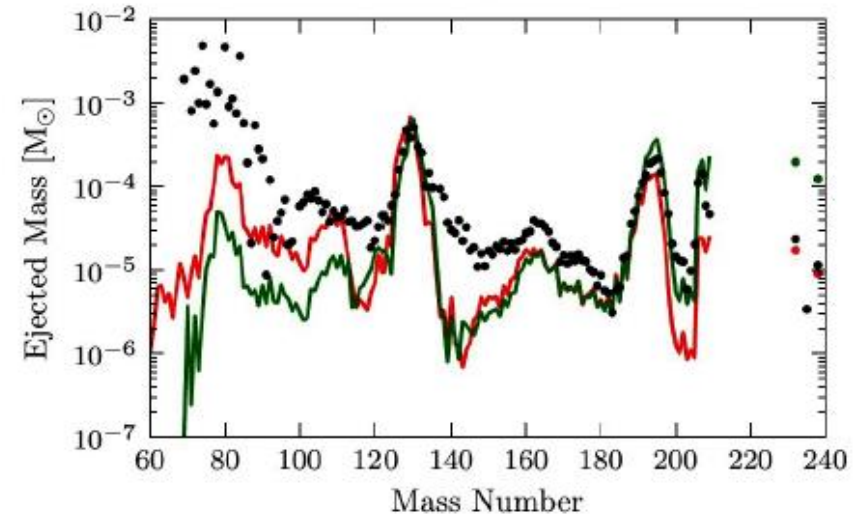
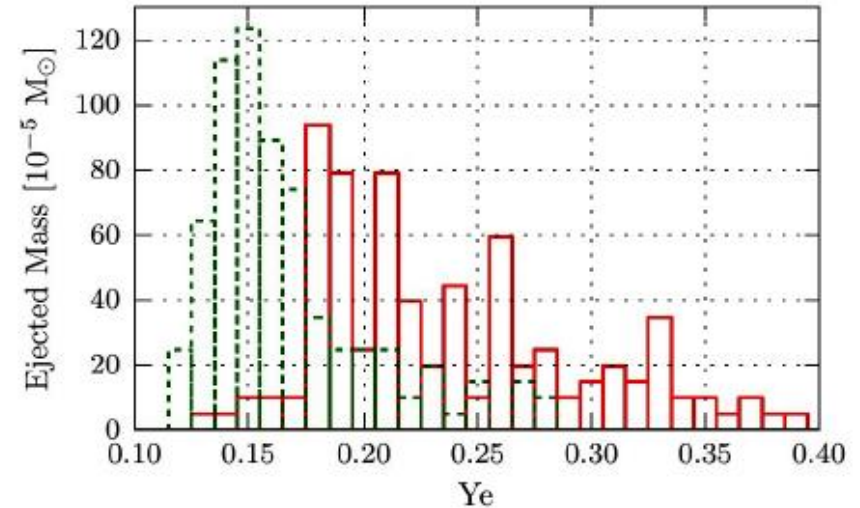
Typical spectra of Stripped-envelope core-collapse SNe



Jet-Supernova Models as r-process Sites?

- MHD-driven polar “jets” could sweep out n-rich matter.
- Requires extremely fast matter ejection, extremely rapid rotation and extremely strong magnetic fields in pre-collapse stellar cores.
- Should be very rare event; maybe 1 of 1000 stellar core collapses?

Winteler et al., ApJL 750 (2012) L22

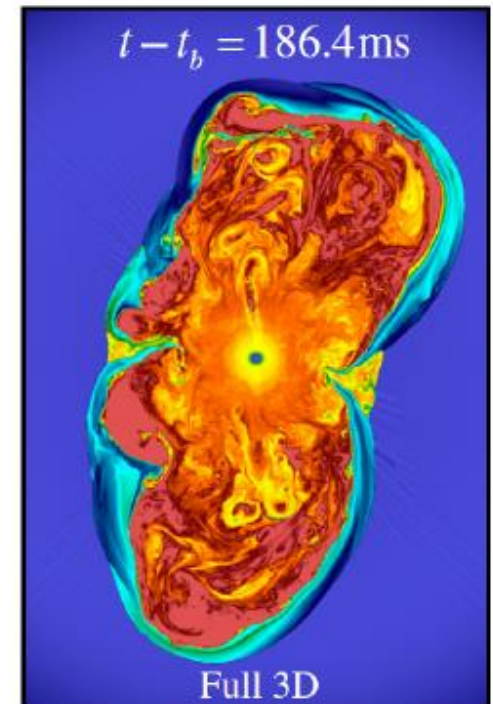
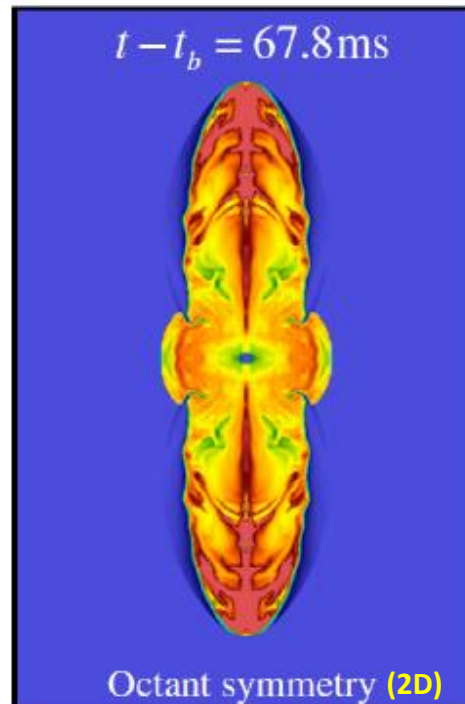


From Th. Janka, 2016

Jet-Supernova Models as r-process Sites?

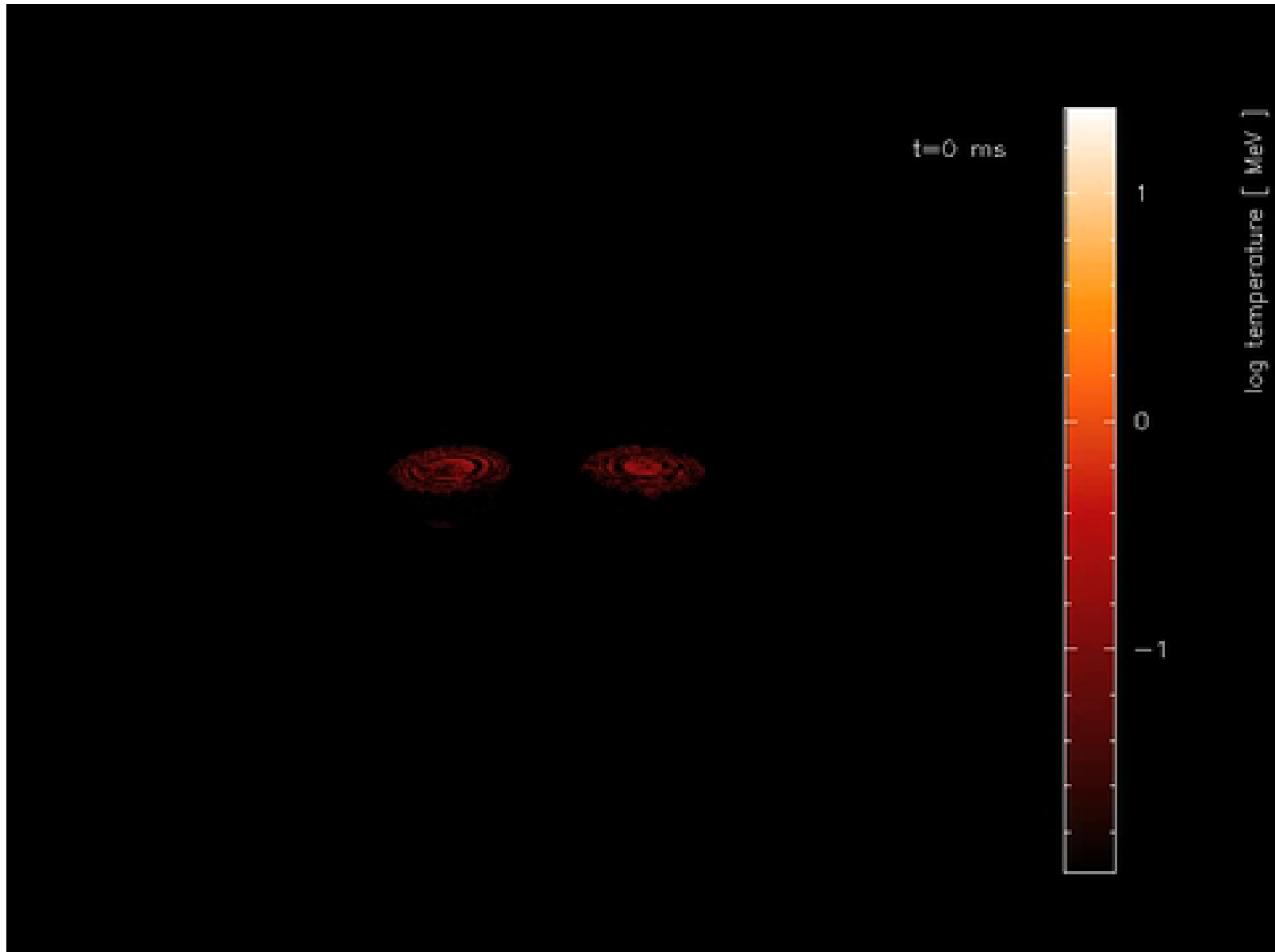
BUT:

- MHD-driven polar “jets” in 3D develop kink instability.
- Assumed initial conditions not supported by stellar pre-collapse models.
- Dynamical scenario does not provide environment for robust r-process.

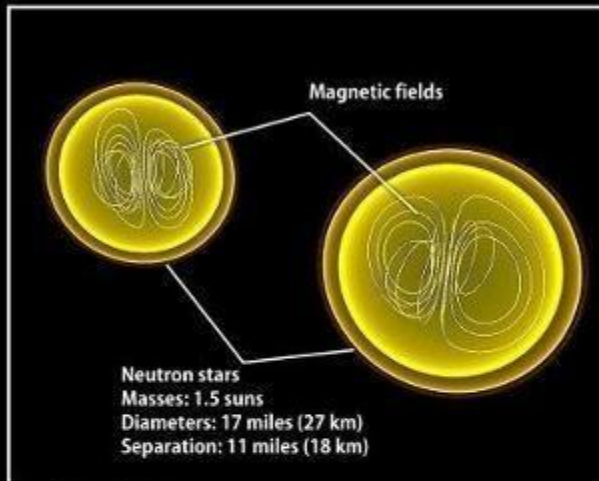


From Th. Janka, 2016

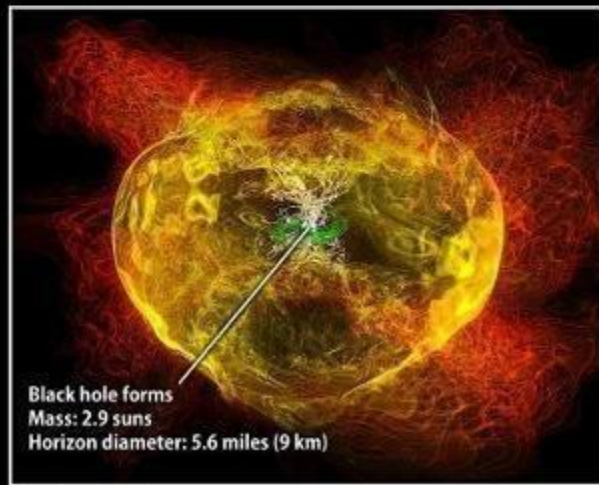
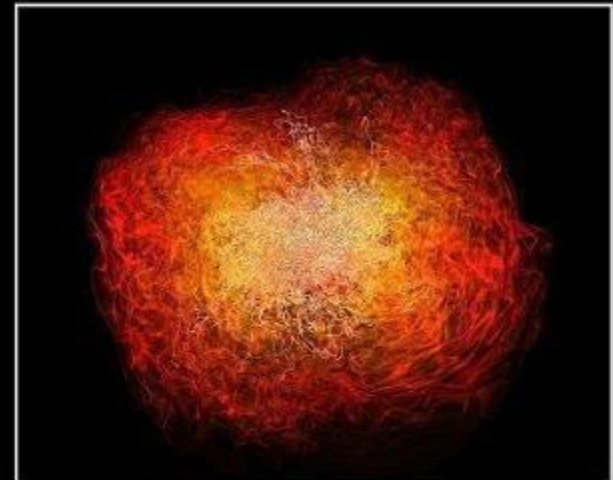
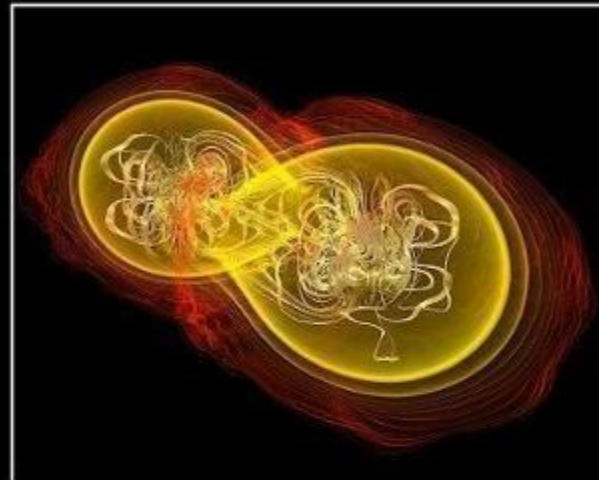
Simulation of the merger of a 1.3 and a 1.4 Msun neutron stars
(S. Rosswog, <http://compact-merger.astro.su.se>)



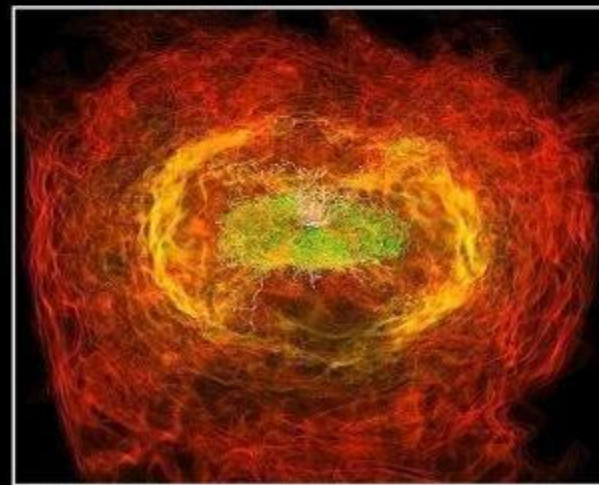
Crashing neutron stars can make gamma-ray burst jets



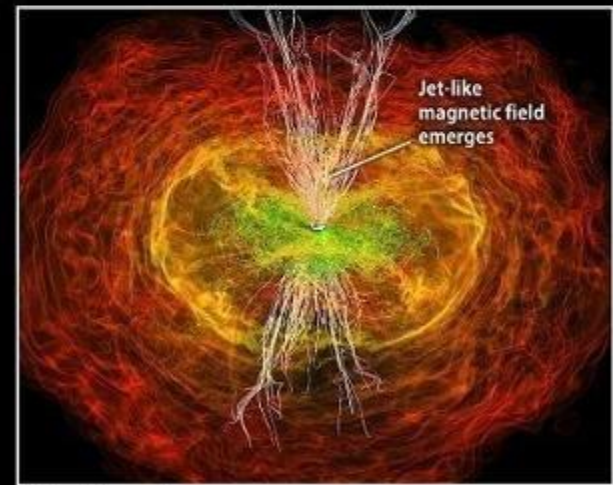
Simulation begins



15.3 milliseconds

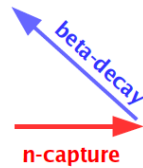
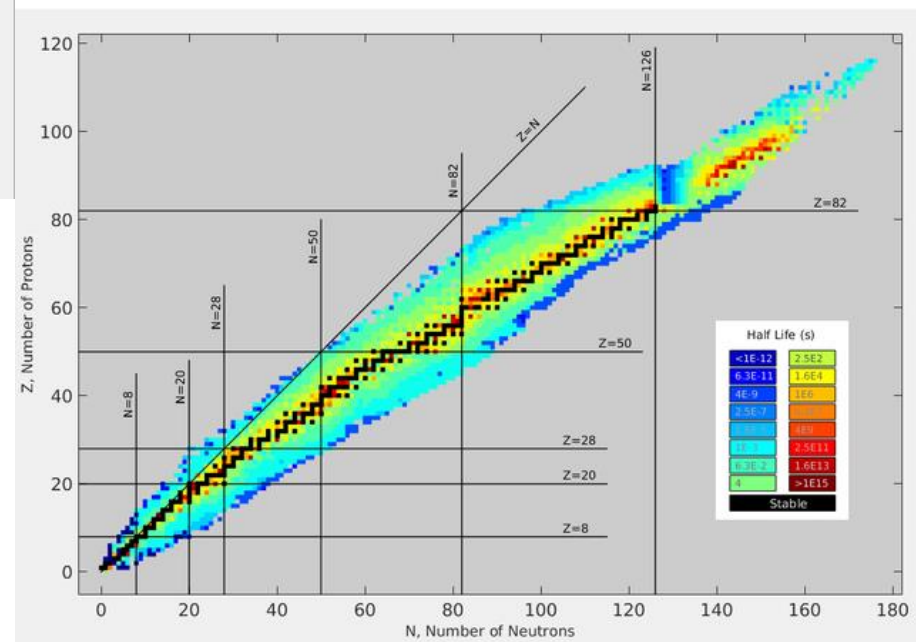
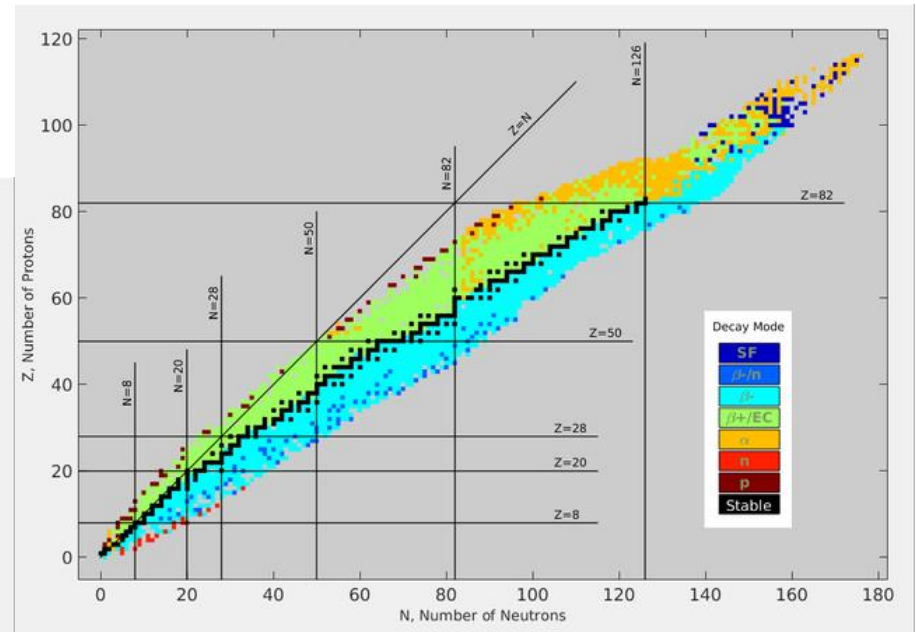
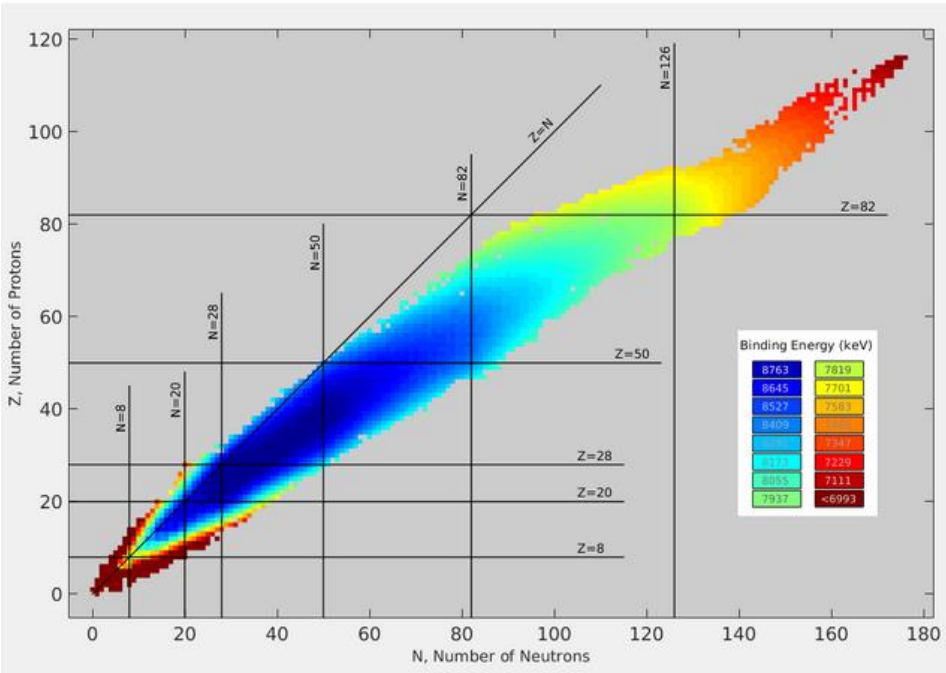


21.2 milliseconds



26.5 milliseconds

The valley of stability



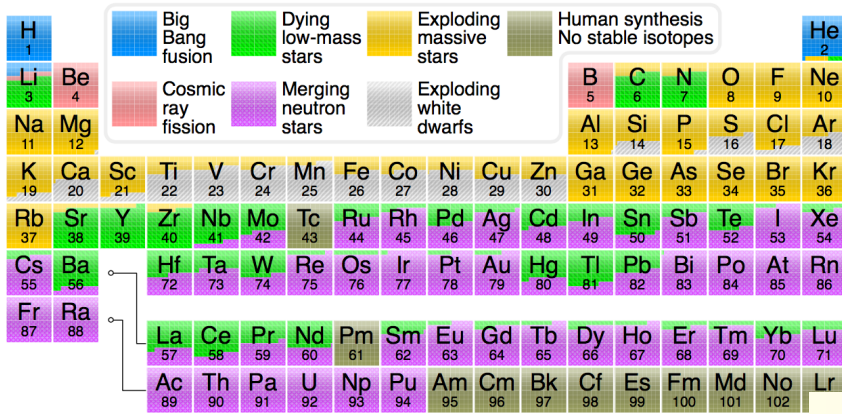
https://en.wikipedia.org/wiki/Valley_of_stability

See also:

<https://www.youtube.com/watch?v=BaUjNiJYgO4>

https://www.youtube.com/watch?v=UTOp_2ZVZmM&t=192s

Periodic table

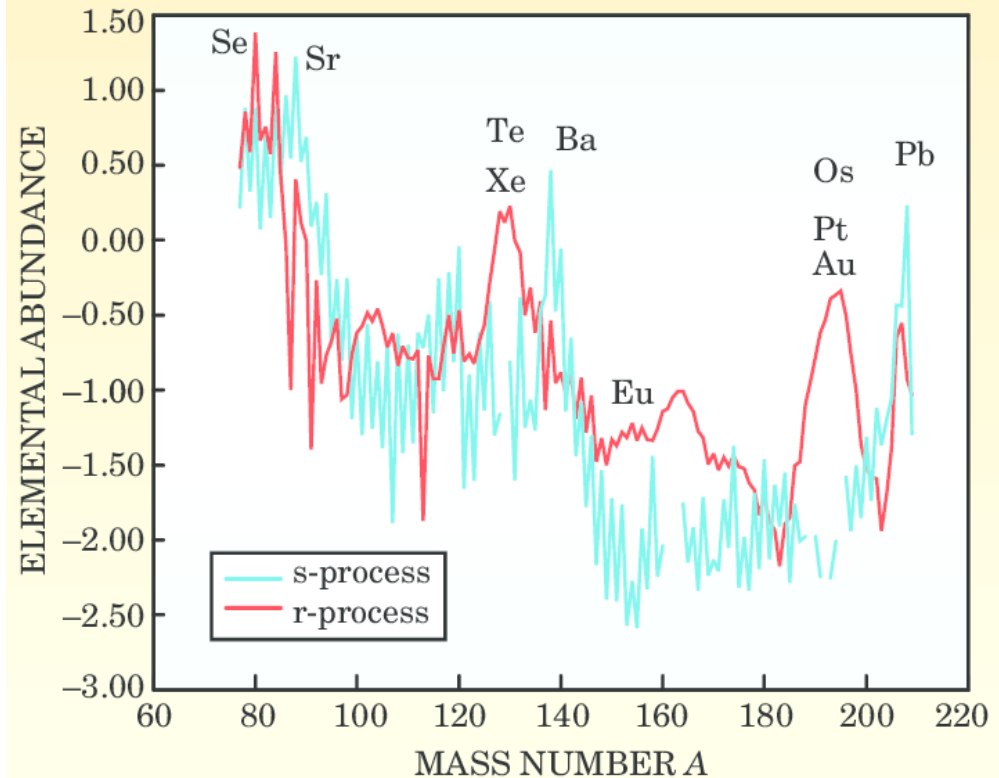


Heavy elements: s- and r-process nucleosynthesis

Solar system abundances of heavy elements produced by r-process and s-process neutron capture.

The s-process (~1000 years) takes place in red-giants cores (AGB phase)

The r-process (seconds or less) takes place in binary NS mergers



Short GRB130603B ($z = 0.356$)

Kilonova:

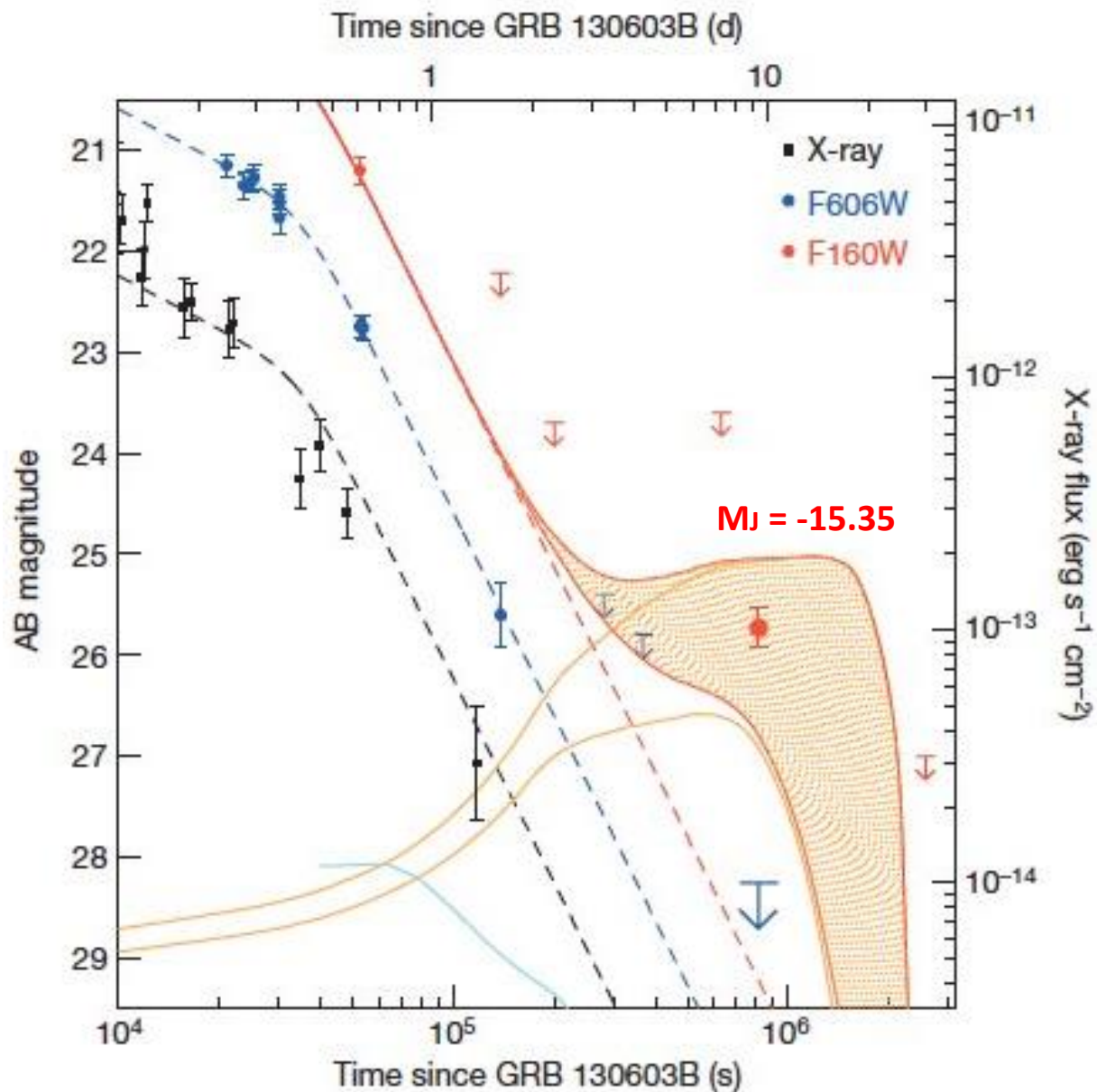
Ejection of r-process material
from a NS merger (0.01-0.1 Mo)
(Barnes & Kasen 2013)

$M_H \approx -15$

$M_R \approx -13$

Tanvir et al. 2013;

Berger et al. 2013



Conclusions

Gravitational waves are a spacetime perturbation, described by a 4-dimensional tensor, produced by massive asymmetric sources.

First gravitational radiation detection (~ 100 Hz) in 2015 lagged theoretical prediction by one century, but results were spectacular: a binary system of black holes of 30 solar masses each merges and produces a nearly 60 solar mass black hole. The mass difference, 3 solar masses, is radiated away as GWs. However, no electromagnetic signal is detected, and none is expected (in principle).

The search for multiwavelength electromagnetic signal accompanying GWs from a binary neutron star merger has manifold implications: compact stellar object physics (NS EoS), binary merger dynamics, GRB physics, nucleosynthesis, MeV neutrino physics.

In particular, the optical/infrared counterpart of a binary NS merger is a radioactive source (kilonova), analogous to supernova.

Nucleosynthesis of binary NS merger points to dominance of r-process, as opposed to s-process: neutron fluxes are rapid and abundant.